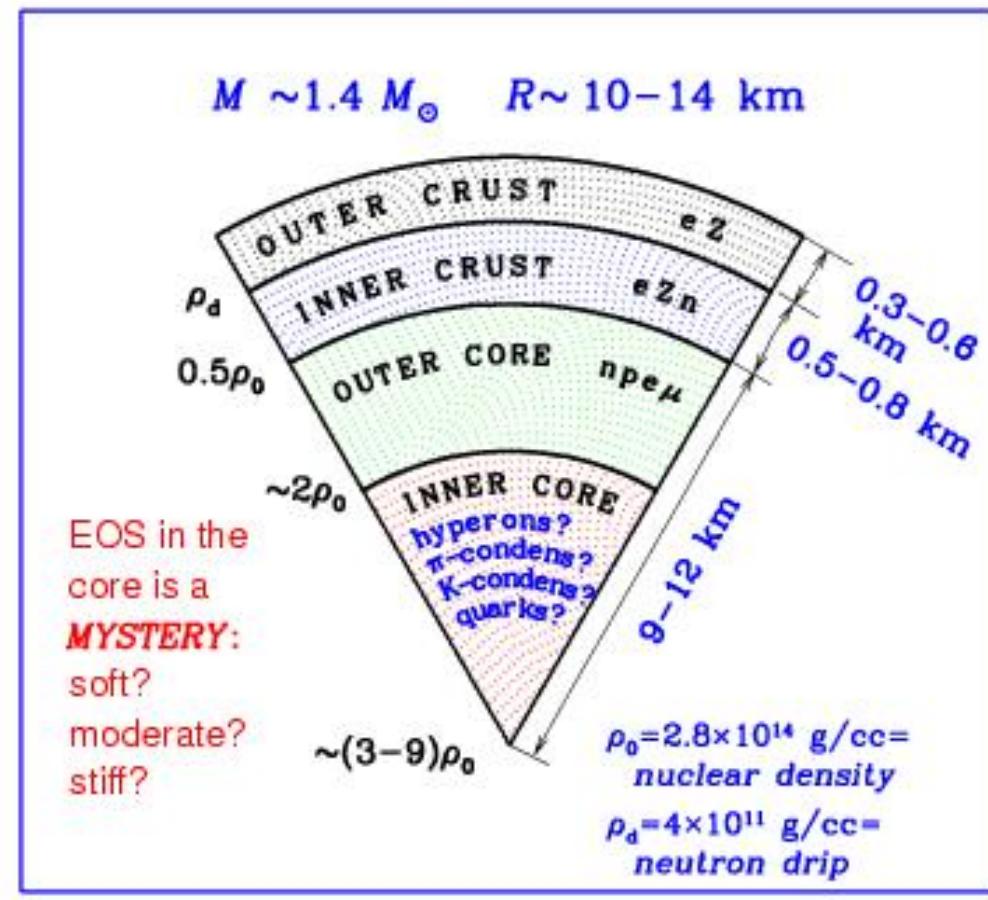
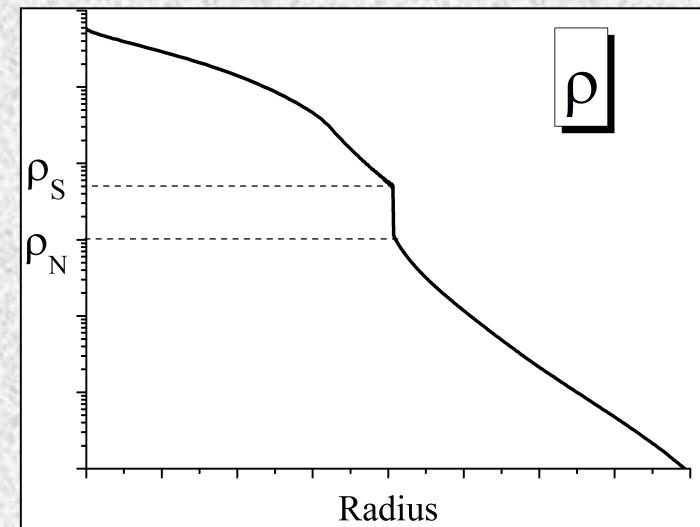
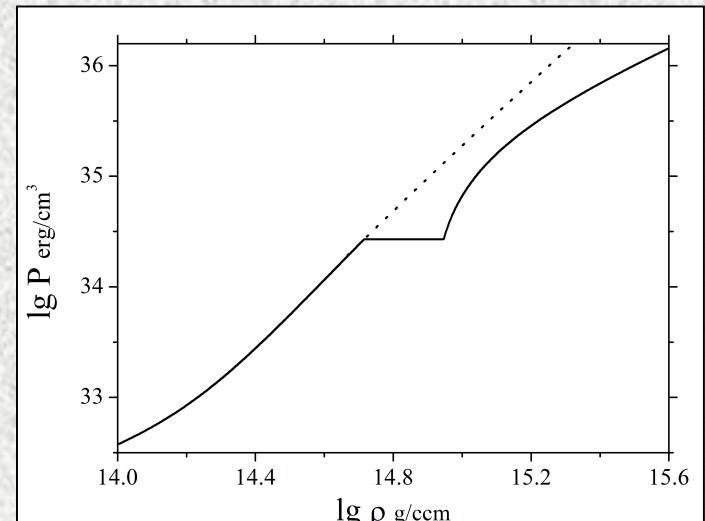


On an amazing feature of the mass-radius relation for superdense hybrid stars

A.V. Yudin, T.L. Razinkova, D.K. Nadyozhin, A.D. Dolgov

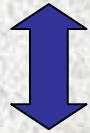
Possible phase transitions in matter at high density



Mass-radius diagram for hybrid stars

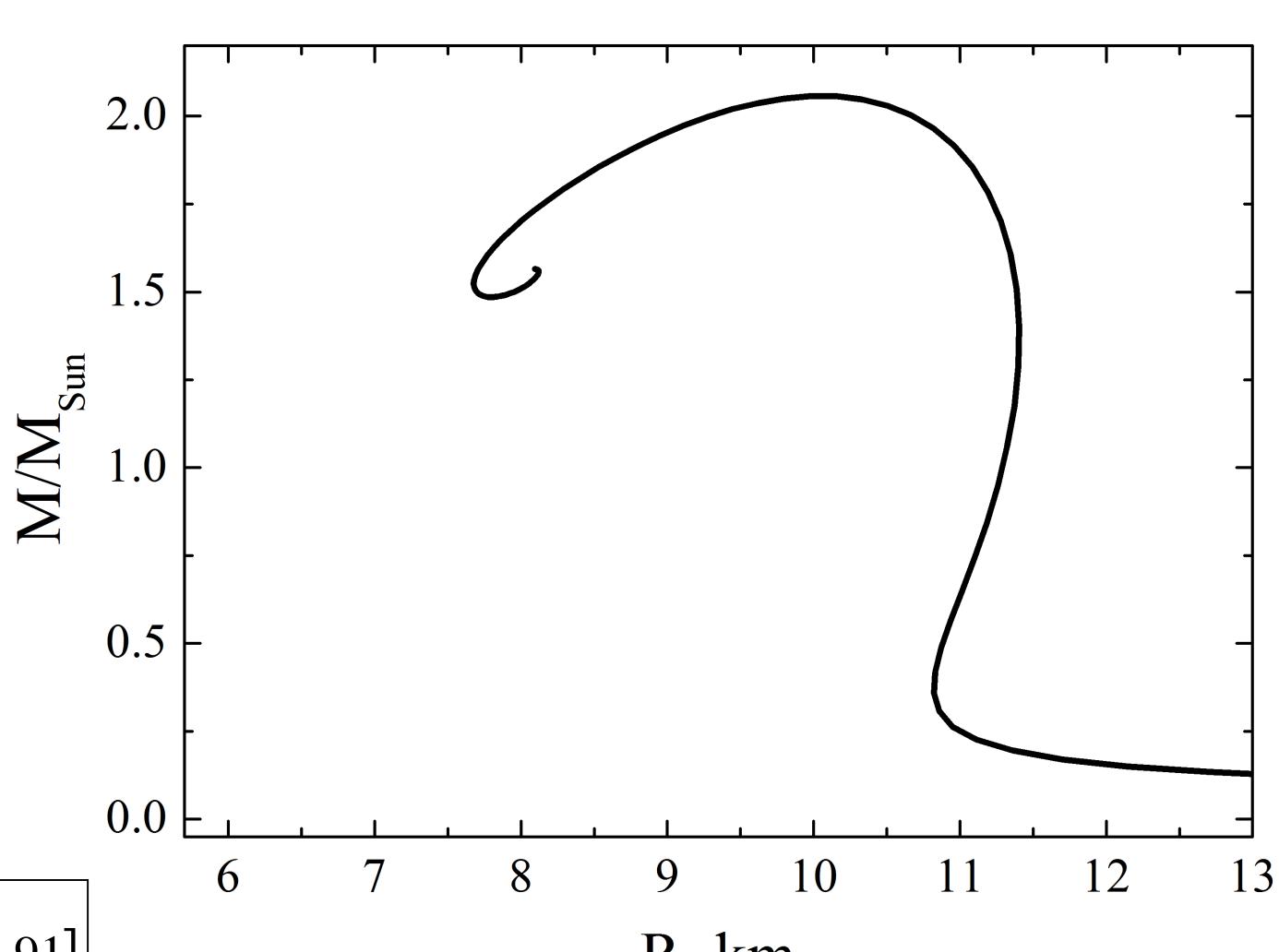
Phase transition:

Nuclear Matter



Quark Matter

$$P_q = \frac{1}{3}(E - 4B)$$



$$\frac{\rho_{pt}}{\rho_n} \approx -3 + \ln[B - 91]$$

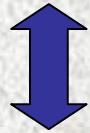
Values of B are in MeV/fm³

EOS for pure nuclear matter is from Douchin & Haensel, Astron. Astrophys., 380 151-167 (2001)

Mass-radius diagram for hybrid stars

Phase transition:

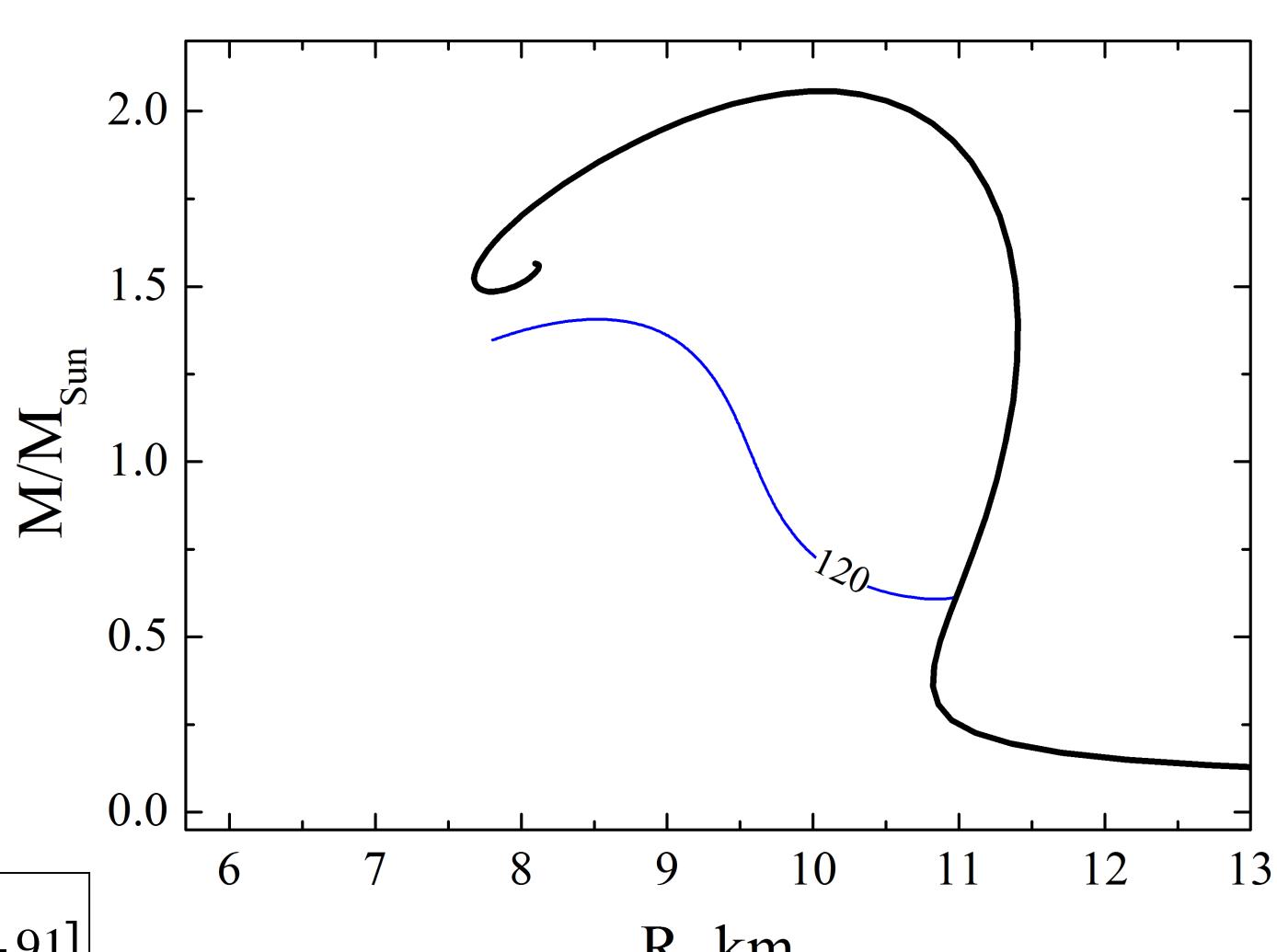
Nuclear Matter



Quark Matter

$$P_q = \frac{1}{3}(E - 4B)$$

$$\frac{\rho_{pt}}{\rho_n} \approx -3 + \ln[B - 91]$$



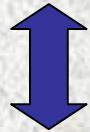
Values of B are in MeV/fm³

EOS for pure nuclear matter is from Douchin & Haensel, Astron. Astrophys., 380 151-167 (2001)

Mass-radius diagram for hybrid stars

Phase transition:

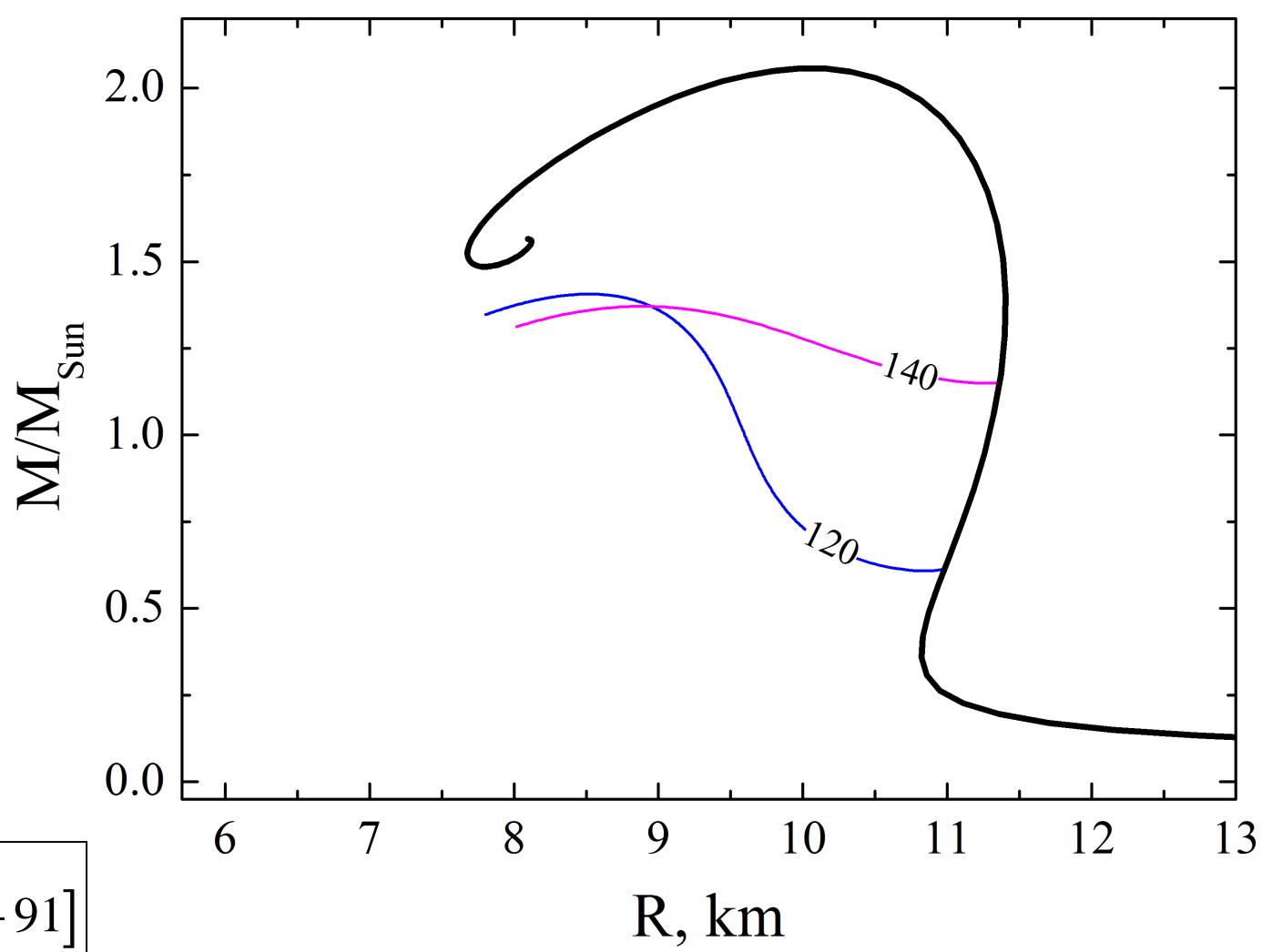
Nuclear Matter



Quark Matter

$$P_q = \frac{1}{3}(E - 4B)$$

$$\frac{\rho_{pt}}{\rho_n} \approx -3 + \ln[B - 91]$$



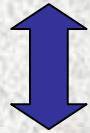
Values of B are in MeV/fm³

EOS for pure nuclear matter is from Douchin & Haensel, Astron. Astrophys., 380 151-167 (2001)

Mass-radius diagram for hybrid stars

Phase transition:

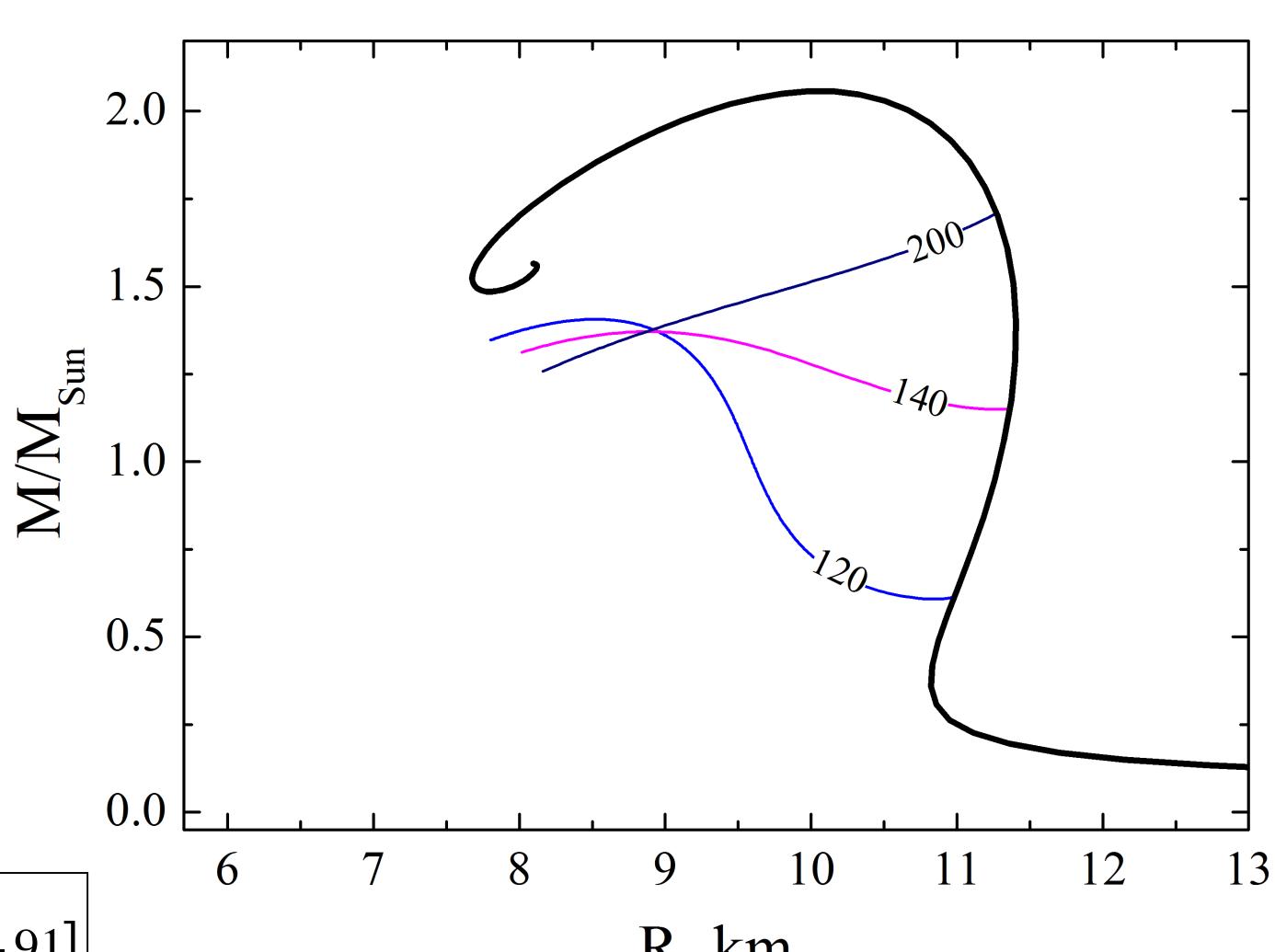
Nuclear Matter



Quark Matter

$$P_q = \frac{1}{3}(E - 4B)$$

$$\frac{\rho_{pt}}{\rho_n} \approx -3 + \ln[B - 91]$$



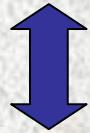
Values of B are in MeV/fm³

EOS for pure nuclear matter is from Douchin & Haensel, Astron. Astrophys., 380 151-167 (2001)

Mass-radius diagram for hybrid stars

Phase transition:

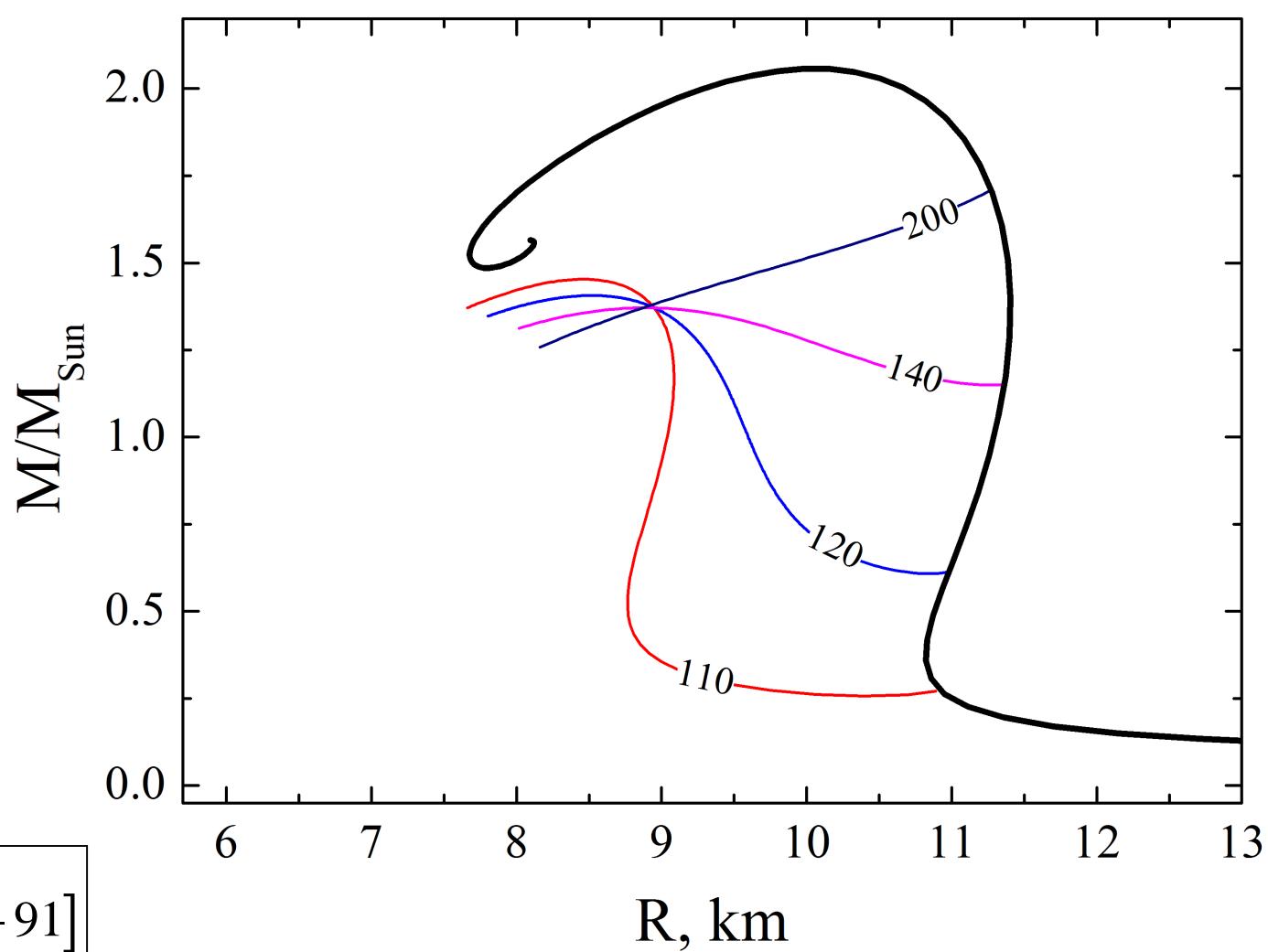
Nuclear Matter



Quark Matter

$$P_q = \frac{1}{3}(E - 4B)$$

$$\frac{\rho_{pt}}{\rho_n} \approx -3 + \ln[B - 91]$$



Values of B are in MeV/fm³

EOS for pure nuclear matter is from Douchin & Haensel, Astron. Astrophys., 380 151-167 (2001)

Mass-radius diagram for hybrid stars

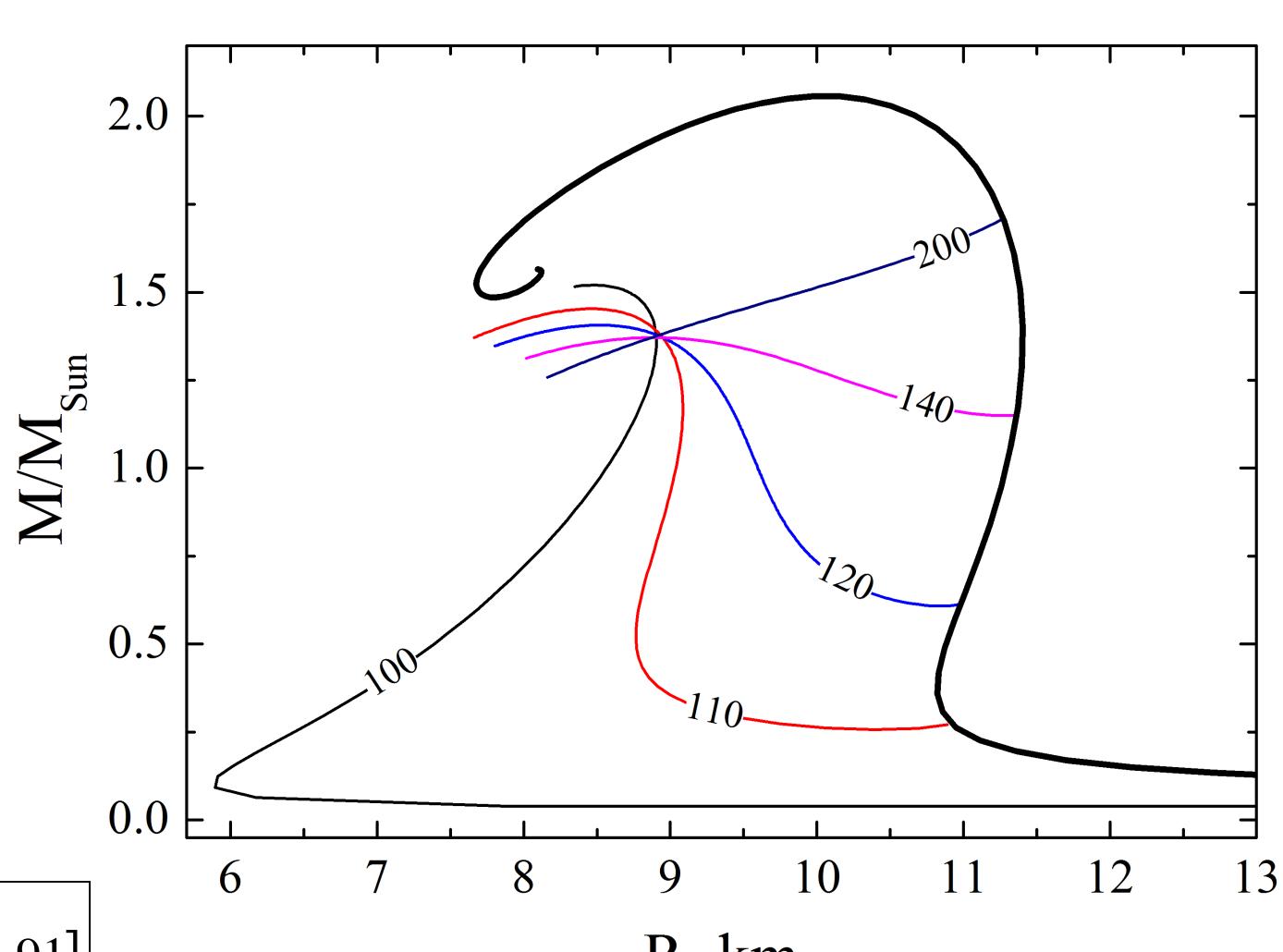
Phase transition:

Nuclear Matter



Quark Matter

$$P_q = \frac{1}{3}(E - 4B)$$



$$\frac{\rho_{pt}}{\rho_n} \approx -3 + \ln[B - 91]$$

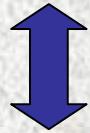
Values of B are in MeV/fm³

EOS for pure nuclear matter is from Douchin & Haensel, Astron. Astrophys., 380 151-167 (2001)

Mass-radius diagram for hybrid stars

Phase transition:

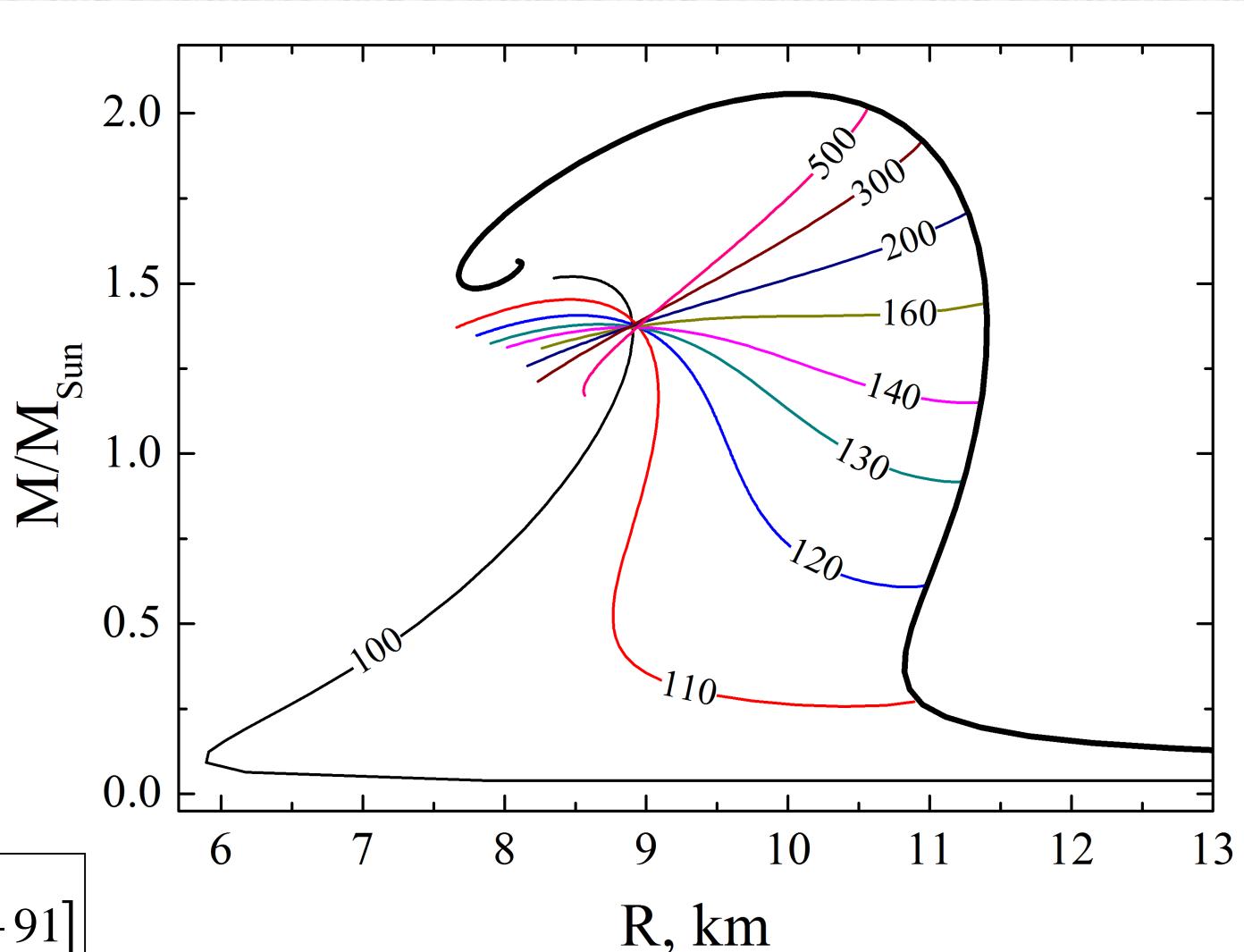
Nuclear Matter



Quark Matter

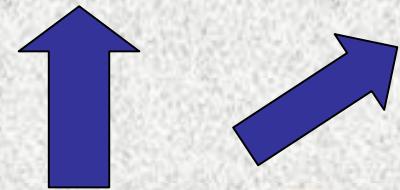
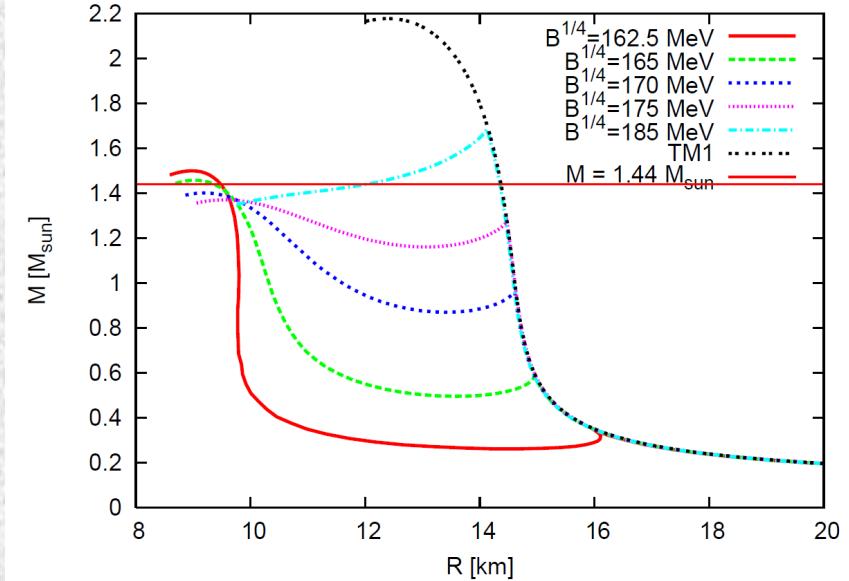
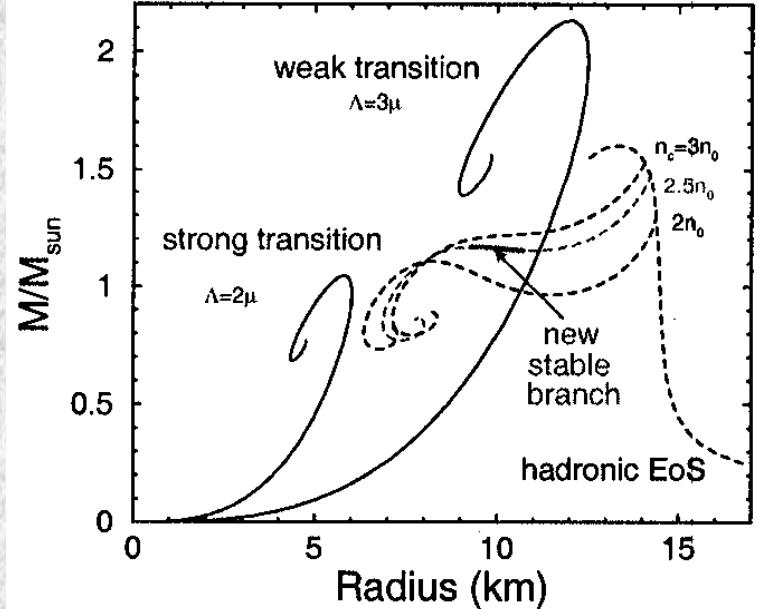
$$P_q = \frac{1}{3}(E - 4B)$$

$$\frac{\rho_{pt}}{\rho_n} \approx -3 + \ln[B - 91]$$

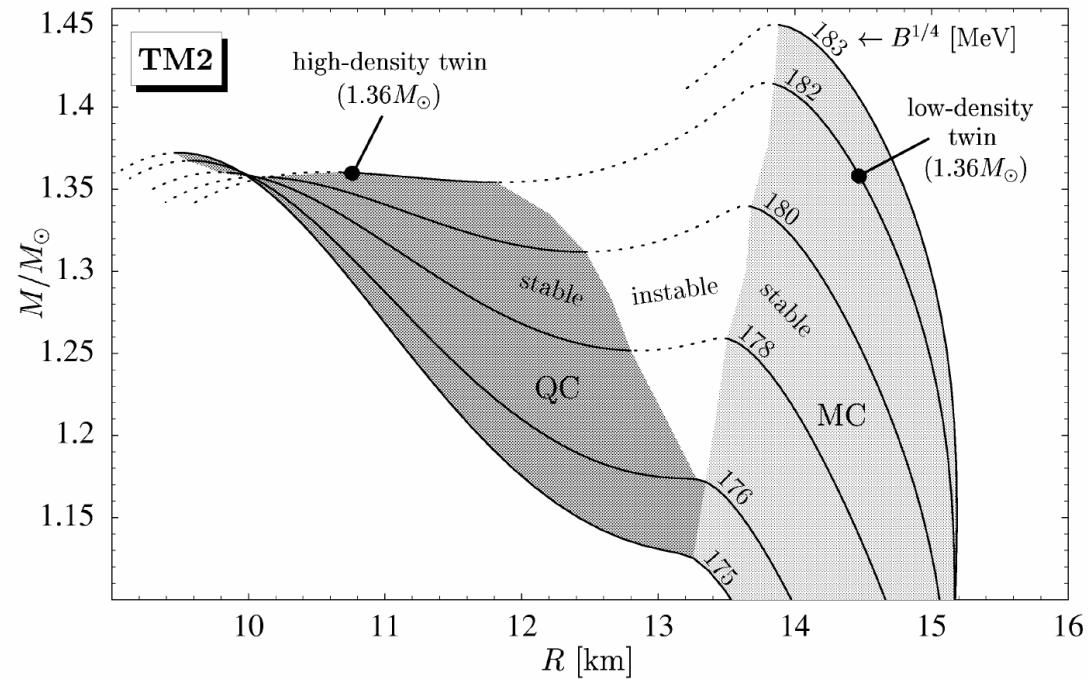


Values of B are in MeV/fm³

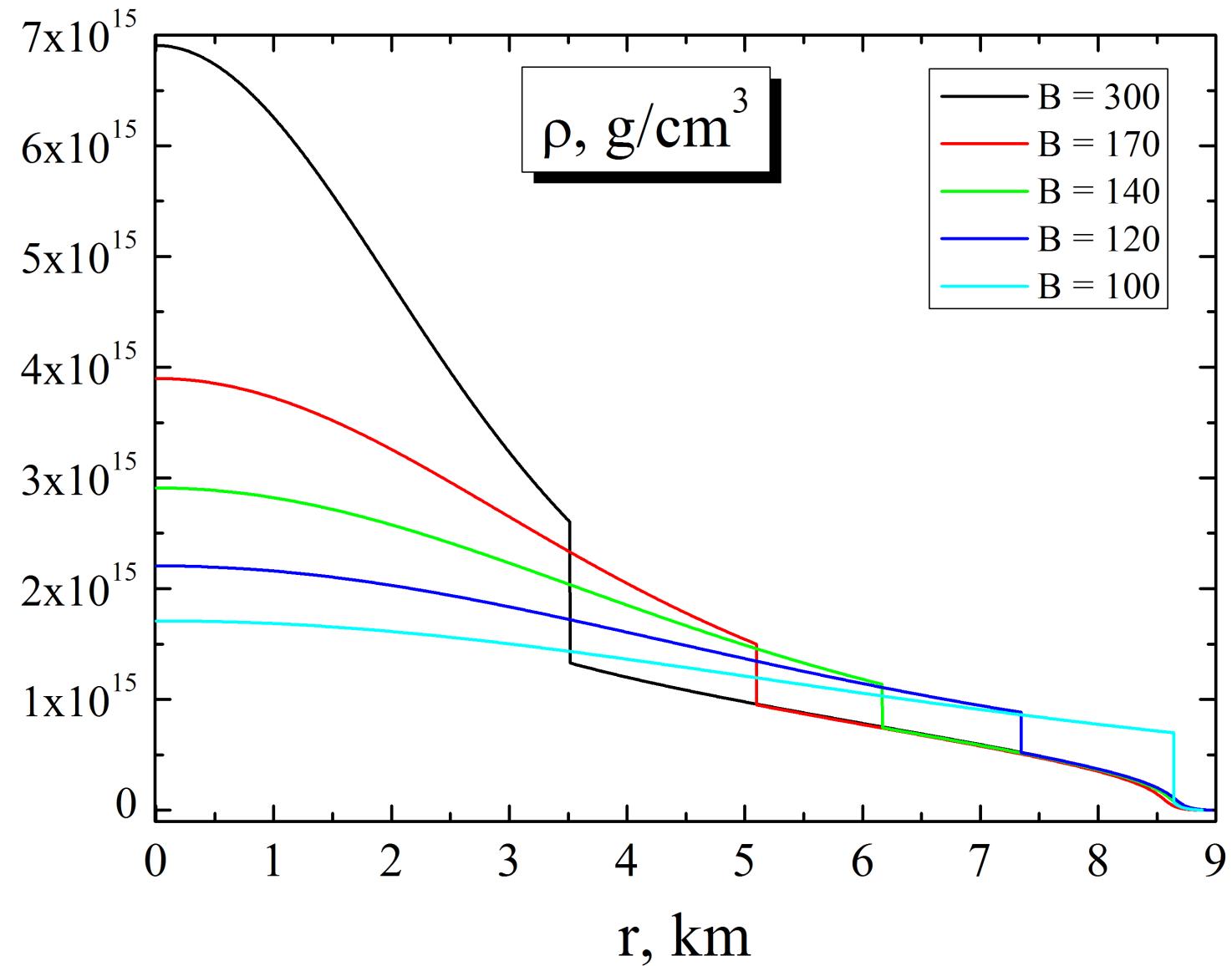
EOS for pure nuclear matter is from Douchin & Haensel, Astron. Astrophys., 380 151-167 (2001)



The are other examples
of such M-R diagram
in the literature!



What is going on inside the star at the Magic Point?



OVT-equations of equilibrium

$$\begin{cases} \frac{dP}{dr} = -\frac{G(P+E)\left(m + \frac{4\pi r^3}{c^2}P\right)}{c^2 r \left(r - \frac{2Gm}{c^2}\right)}, \\ \frac{dm}{dr} = 4\pi r^2 \frac{E}{c^2}. \end{cases}$$



$$\Delta P = \left(\frac{\partial P}{\partial r} \right)_1 \delta r = \frac{1}{\lambda} \left(\frac{\partial P}{\partial r} \right)_2 \delta r.$$

$$\Delta m = 4\pi r_*^2 \frac{E_1}{c^2} \delta r = 4\pi r_*^2 \frac{E_2}{c^2} \left[1 - (\lambda - 1) \frac{P_2}{E_2} \right] \frac{\delta r}{\lambda}.$$

$$\begin{cases} \Delta P = \left(\frac{dP}{dr} \right)_2 \delta r + \left(\frac{dP}{dP_c} \right)_{r,\zeta} \delta P_c + \left(\frac{dP}{d\zeta} \right)_{r,P_c} \delta \zeta, \\ \Delta m = \left(\frac{dm}{dr} \right)_2 \delta r + \left(\frac{dm}{dP_c} \right)_{r,\zeta} \delta P_c + \left(\frac{dm}{d\zeta} \right)_{r,P_c} \delta \zeta. \end{cases}$$

Phase equilibrium

$$\begin{cases} P_{pt} = P_1(n_1) = P_2(n_2) \\ \mu_{pt} = \frac{P_1 + E_1}{n_1} = \frac{P_2 + E_2}{n_2} \end{cases}$$

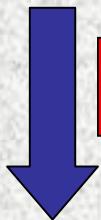
$$E_2 = E_2(n_2, \zeta), \quad \zeta \equiv B \text{ for quark matter}$$



$$\Delta P = \frac{\delta \zeta}{\lambda - 1} \left(\frac{\partial E_2}{\partial \zeta} \right), \quad \lambda \equiv \frac{n_2}{n_1}.$$

The main equation

$$\begin{aligned}\frac{dP}{dr} \left[\frac{\lambda - 1}{\lambda} \right] \delta r &= \left(\frac{\partial E_2}{\partial \zeta} \right) \delta \zeta, \\ -\frac{dP}{dr} \left[\frac{\lambda - 1}{\lambda} \right] \delta r &= \left(\frac{\partial P}{\partial P_c} \right)_{r,\zeta} \delta P_c + \left(\frac{\partial P}{\partial \zeta} \right)_{r,P_c} \delta \zeta, \\ -\frac{dm}{dr} \left[\frac{P_2 + E_2}{E_2} \right] \left[\frac{\lambda - 1}{\lambda} \right] \delta r &= \left(\frac{\partial m}{\partial P_c} \right)_{r,\zeta} \delta P_c + \left(\frac{\partial m}{\partial \zeta} \right)_{r,P_c} \delta \zeta.\end{aligned}$$

 **Det = 0!**

$$\left(\frac{\partial P}{\partial P_c} \right)_{r,\zeta} \left[\frac{dm}{dr} \frac{(P_2 + E_2)}{E_2} \left(\frac{\partial E_2}{\partial \zeta} \right) + \left(\frac{\partial m}{\partial \zeta} \right)_{r,P_c} \right] = \left(\frac{\partial m}{\partial P_c} \right)_{r,\zeta} \left[\left(\frac{\partial E_2}{\partial \zeta} \right) + \left(\frac{\partial P}{\partial \zeta} \right)_{r,P_c} \right].$$

Dimensionless form of the Main Equation

$$P = \alpha(E - E_0), \quad E_0 \equiv 4B, \quad \alpha_q = \frac{1}{3}$$

$$\rho = \frac{E}{E_0}, \quad x = \frac{r}{r_{\text{dim}}}, \quad \mu = \frac{m}{m_{\text{dim}}}$$

$$r_{\text{dim}} = \frac{c^2}{\sqrt{4\pi G E_0}}, \quad m_{\text{dim}} = \frac{c^4}{G \sqrt{4\pi G E_0}}.$$

Equilibrium equations:

$$\begin{cases} \alpha \frac{d\rho}{dx} = -[\rho + \alpha(\rho - 1)] \frac{\mu + \alpha x^3(\rho - 1)}{x(x - 2\mu)}, \\ \frac{d\mu}{dx} = x^2 \rho. \end{cases}$$

$$\left(\frac{\partial \mu}{\partial \rho_c} \right)_x \left[\rho - \frac{\alpha}{1+\alpha} + \frac{x}{2} \frac{d\rho}{dx} \right] \frac{d\rho}{dx} = x^2 \left(\frac{\partial \rho}{\partial \rho_c} \right)_x \left[\rho - \frac{\alpha}{1+\alpha} + \frac{x}{2} \frac{d\rho}{dx} \left(\rho - \frac{\mu}{x^3} \right) \right]$$

Homology-invariant variables

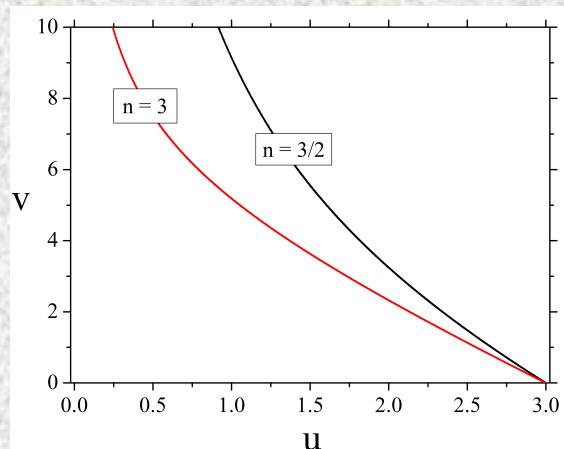
For some specific EOSes TWO TOV-equations of equilibrium can be combined to ONE differential equation with homological variables.

$$\begin{cases} x \frac{du}{dx} = f_u(u, v), \\ x \frac{dv}{dx} = f_v(u, v). \end{cases} \Rightarrow \frac{dv}{du} = \frac{f_v(u, v)}{f_u(u, v)} = f(u, v).$$

Examples:

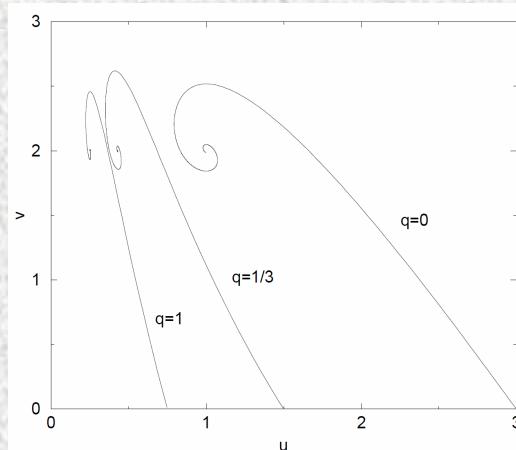
Newtonian limit, polytropes:

$$\begin{cases} u = \frac{d \ln m}{d \ln r} = \frac{4\pi r^3 \rho}{m}, \\ v = -\frac{d \ln P}{d \ln r} = \frac{G m \rho}{r P}. \end{cases} \Rightarrow \frac{u}{v} \frac{dv}{du} = -\frac{u + \frac{v}{1+n} - 1}{u + \frac{nv}{1+n} - 3}$$



GR, P=qE

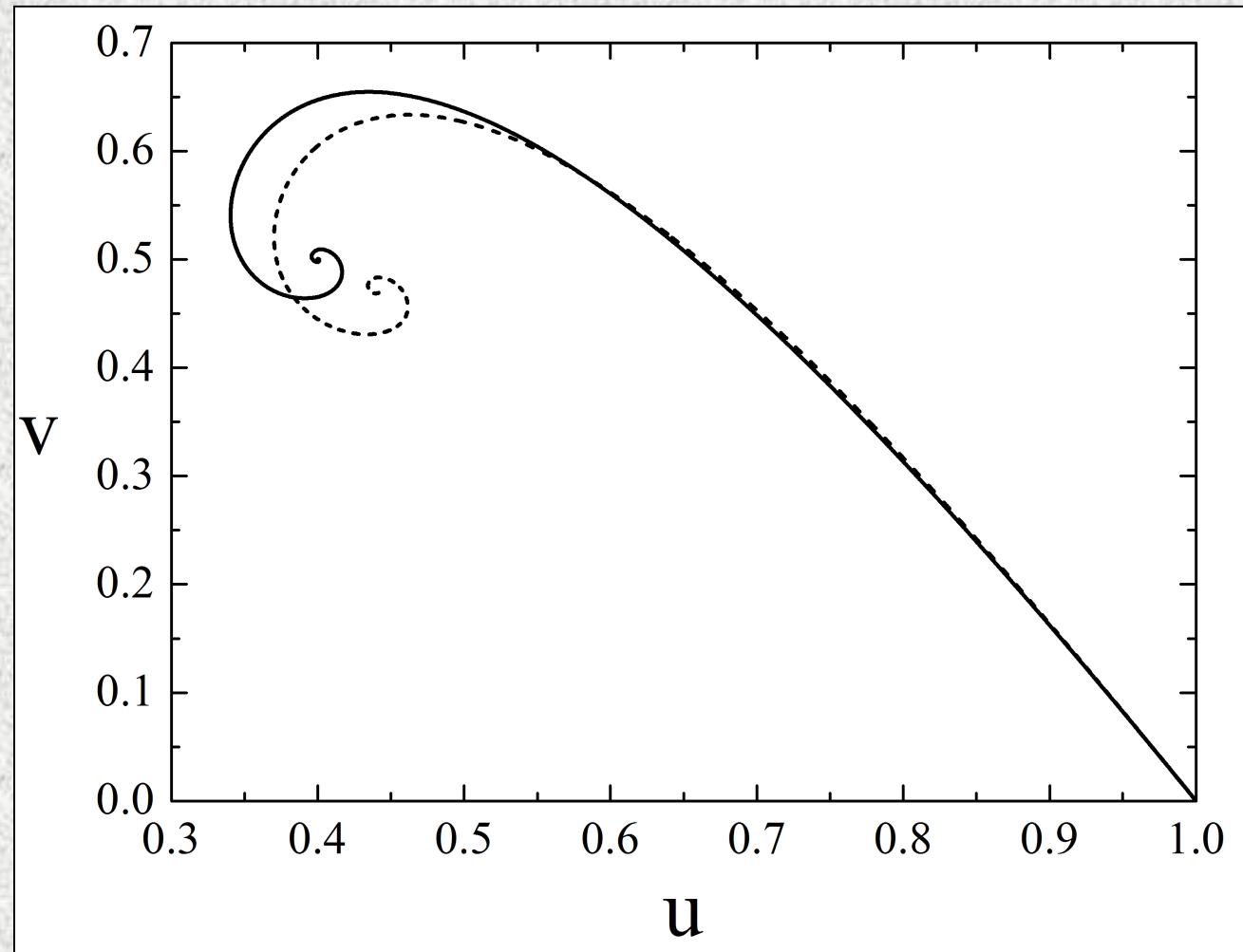
$$\frac{u}{v} \frac{dv}{du} = \frac{-1 - \frac{2q}{1+q}v + (1+3q)u + \frac{q(3+5q)}{1+q}uv + \frac{2q^2(1-q)}{(1+q)^2}uv^2}{3 - \frac{1-q}{1+q}v - (1+3q)u - \frac{q(3+q)}{1+q}uv - \frac{4q^2}{(1+q)^2}uv^2}.$$



From P.-H. Chavanis, Astron. Astrophys., 381, (2002)

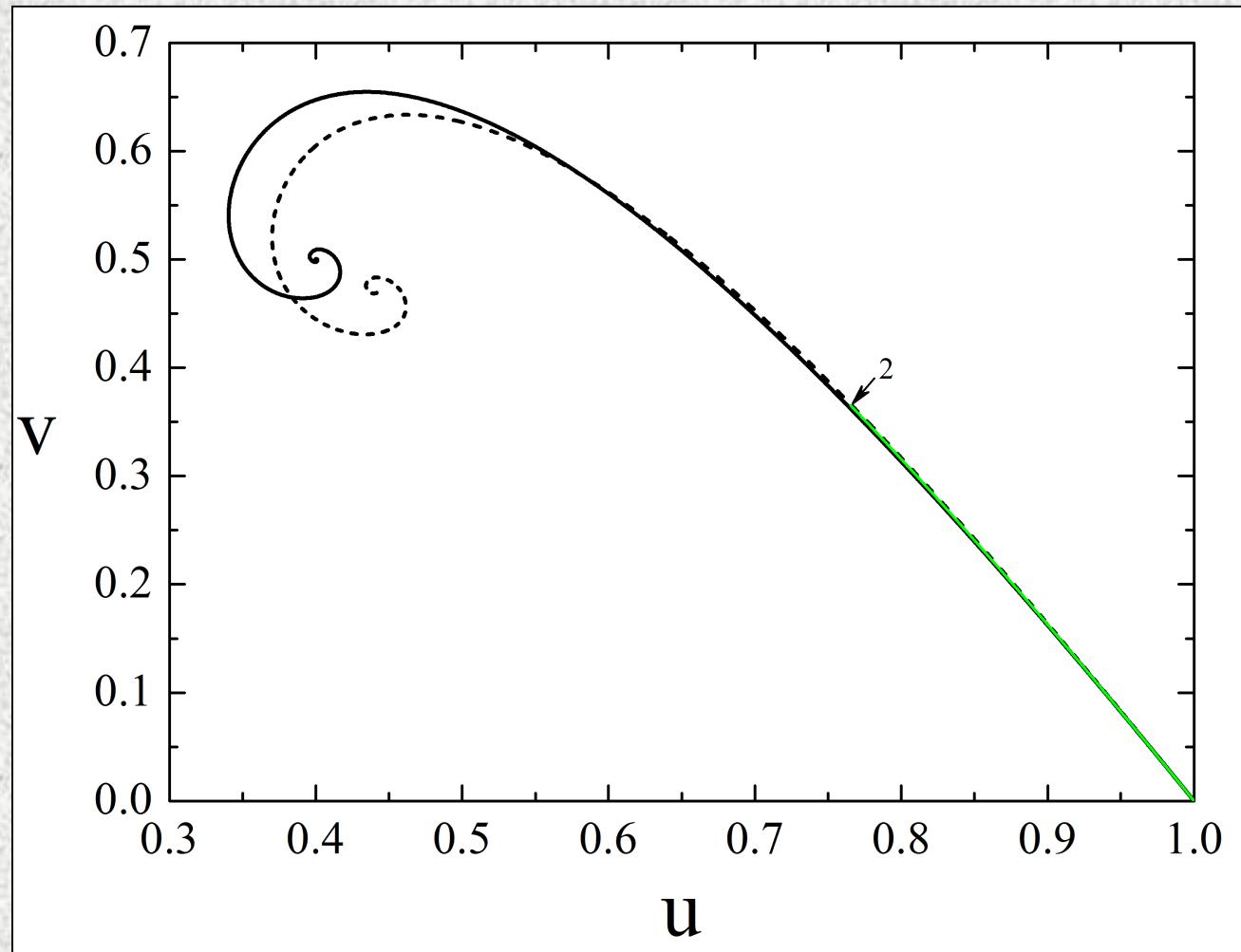
“Almost” homological variables

$$\begin{cases} v = \frac{\mu + \alpha x^3(\rho - 1)}{x - 2\mu}, \\ u = \frac{x^3[\rho + \alpha(\rho - 1)]}{3\mu + \alpha x^3(\rho - 1)}. \end{cases}$$



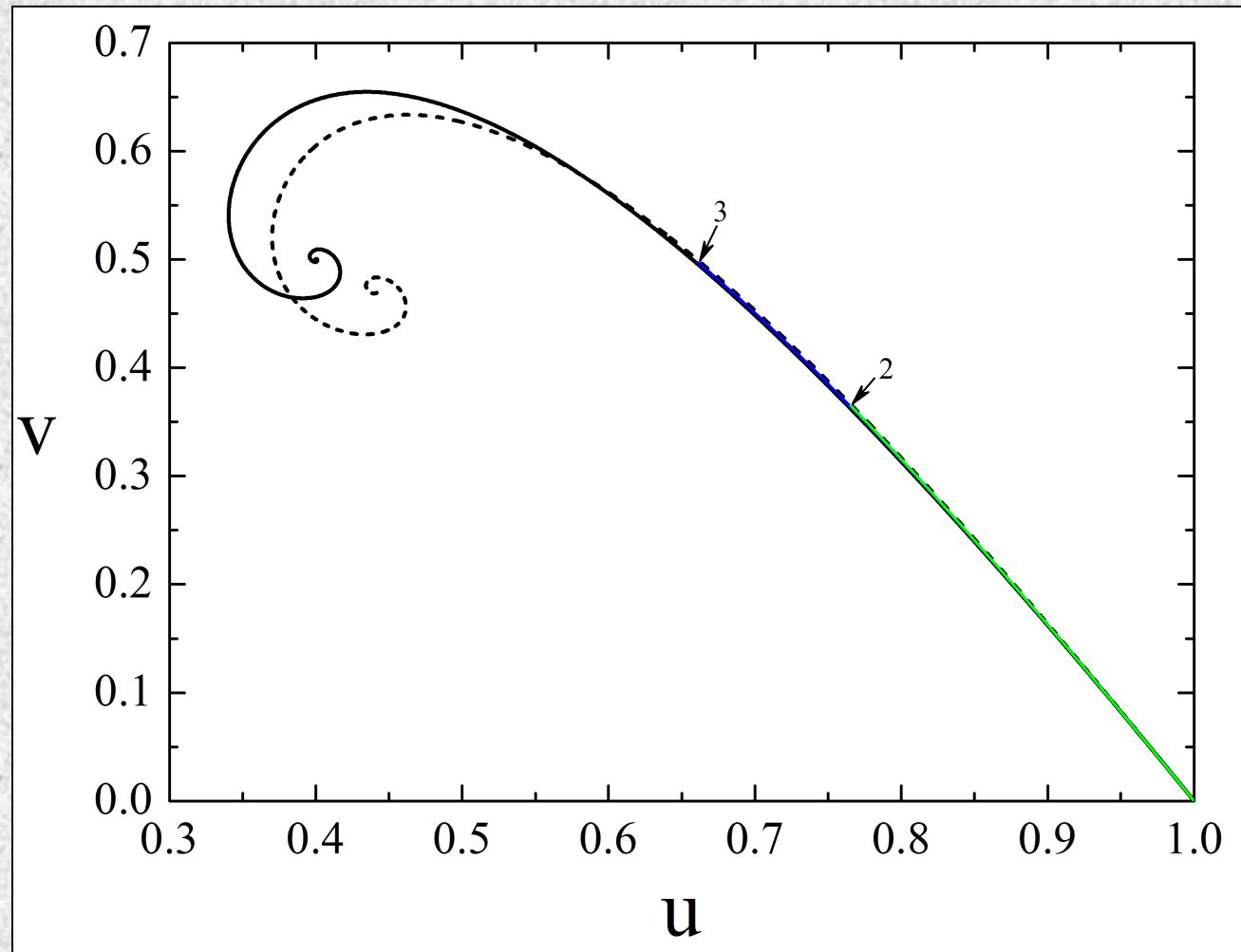
“Almost” homological variables

$$\begin{cases} v = \frac{\mu + \alpha x^3(\rho - 1)}{x - 2\mu}, \\ u = \frac{x^3[\rho + \alpha(\rho - 1)]}{3\mu + \alpha x^3(\rho - 1)}. \end{cases}$$



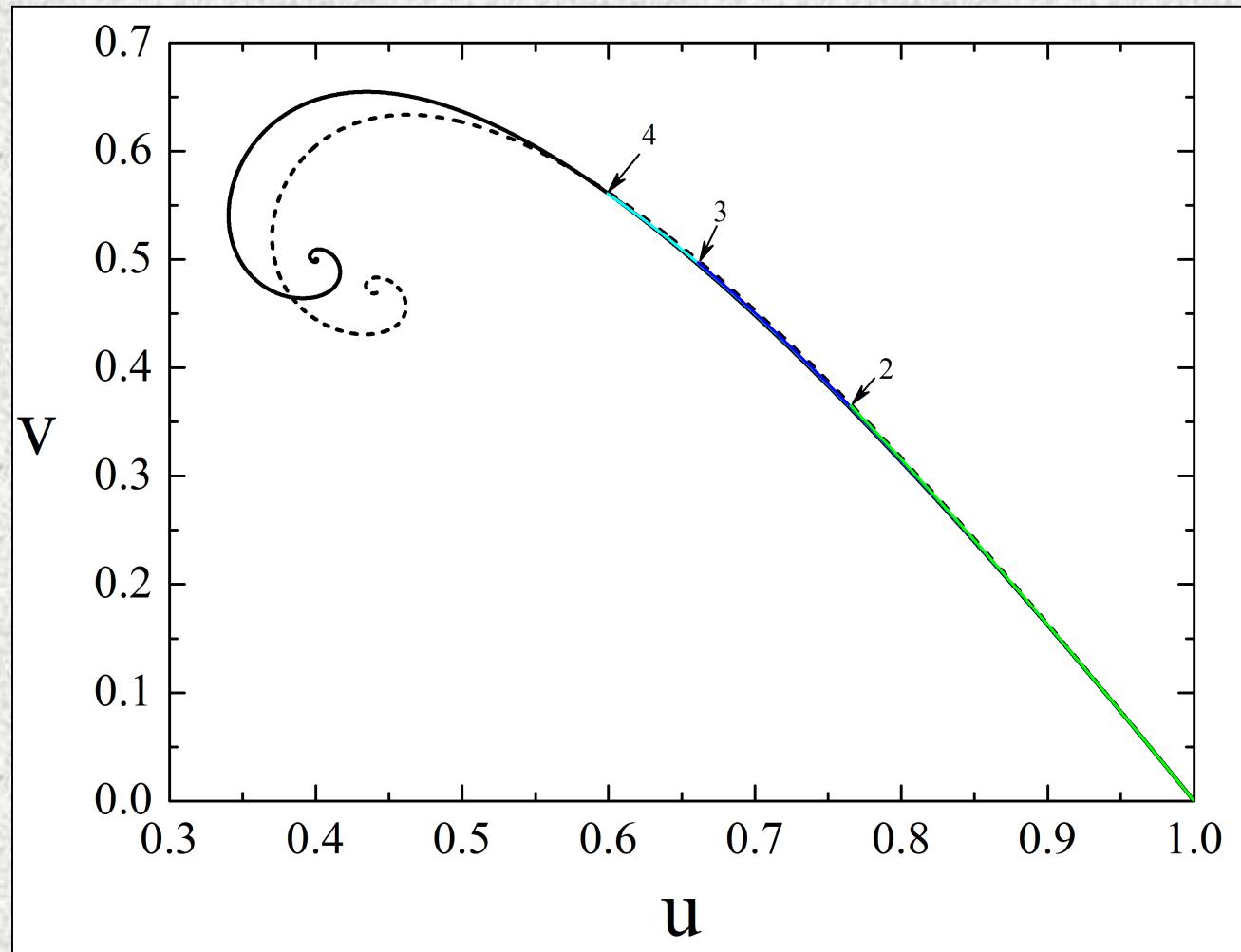
“Almost” homological variables

$$\begin{cases} v = \frac{\mu + \alpha x^3(\rho - 1)}{x - 2\mu}, \\ u = \frac{x^3[\rho + \alpha(\rho - 1)]}{3\mu + \alpha x^3(\rho - 1)}. \end{cases}$$



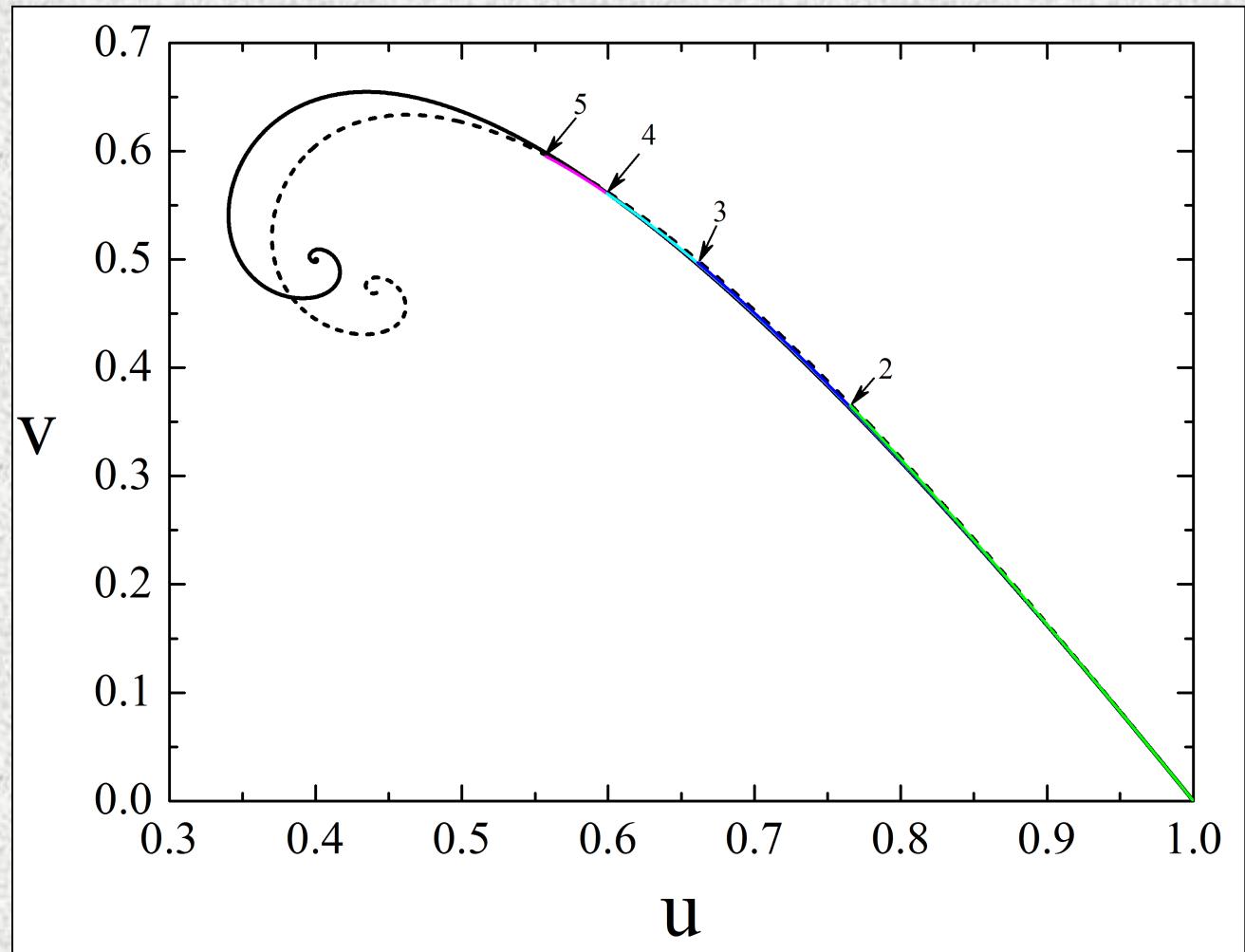
“Almost” homological variables

$$\begin{cases} v = \frac{\mu + \alpha x^3(\rho - 1)}{x - 2\mu}, \\ u = \frac{x^3[\rho + \alpha(\rho - 1)]}{3\mu + \alpha x^3(\rho - 1)}. \end{cases}$$



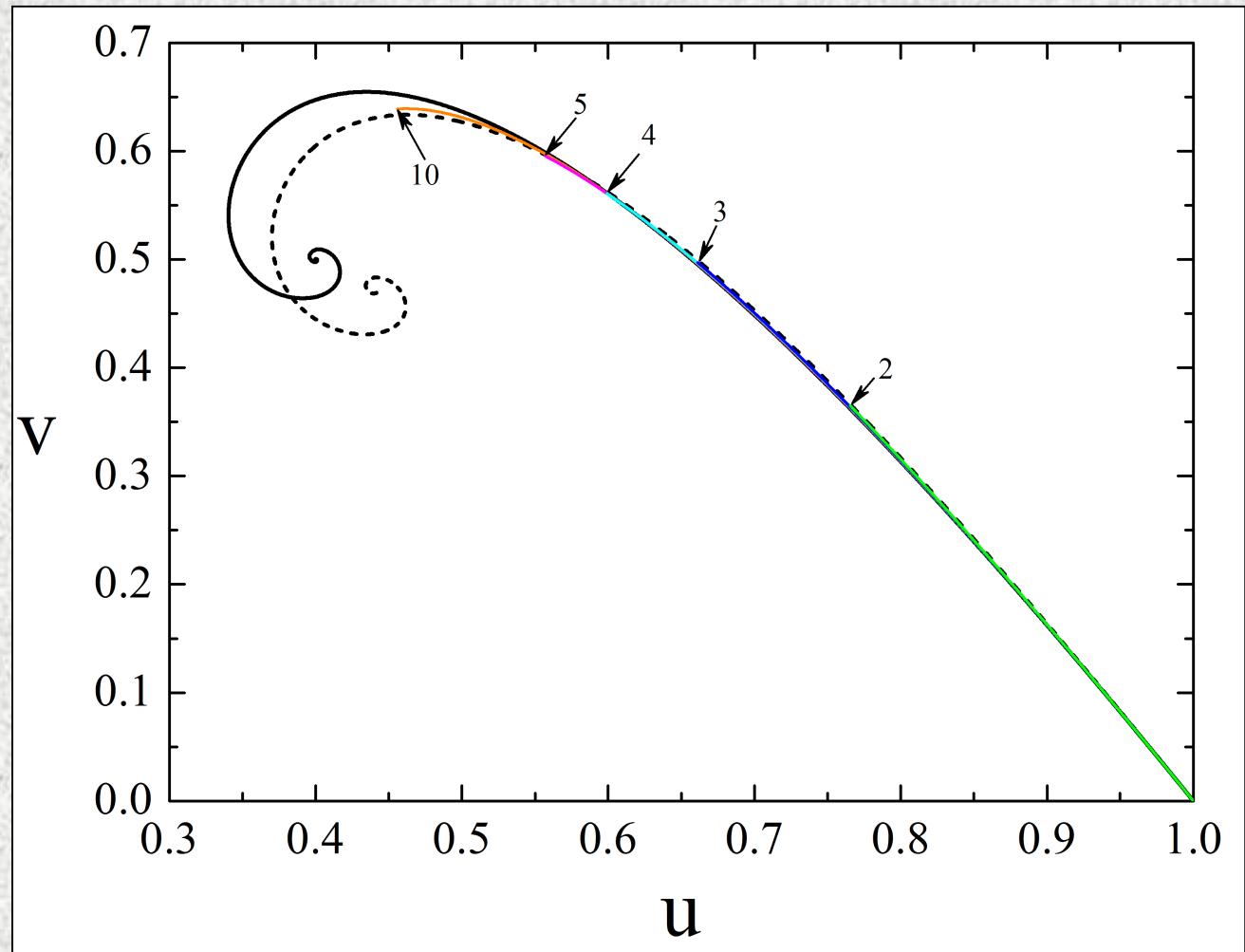
“Almost” homological variables

$$\begin{cases} v = \frac{\mu + \alpha x^3(\rho - 1)}{x - 2\mu}, \\ u = \frac{x^3[\rho + \alpha(\rho - 1)]}{3\mu + \alpha x^3(\rho - 1)}. \end{cases}$$

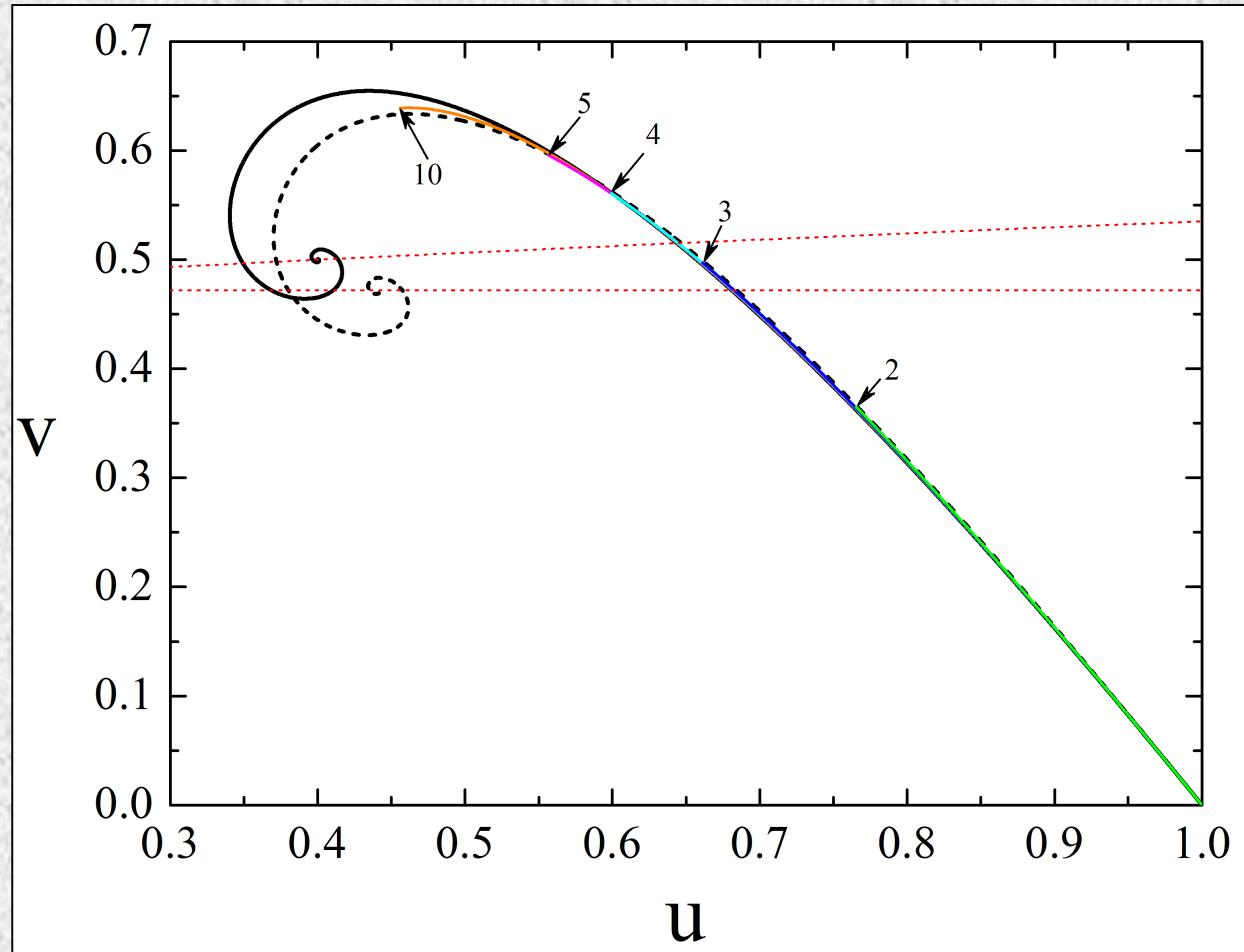


“Almost” homological variables

$$\begin{cases} v = \frac{\mu + \alpha x^3(\rho - 1)}{x - 2\mu}, \\ u = \frac{x^3[\rho + \alpha(\rho - 1)]}{3\mu + \alpha x^3(\rho - 1)}. \end{cases}$$

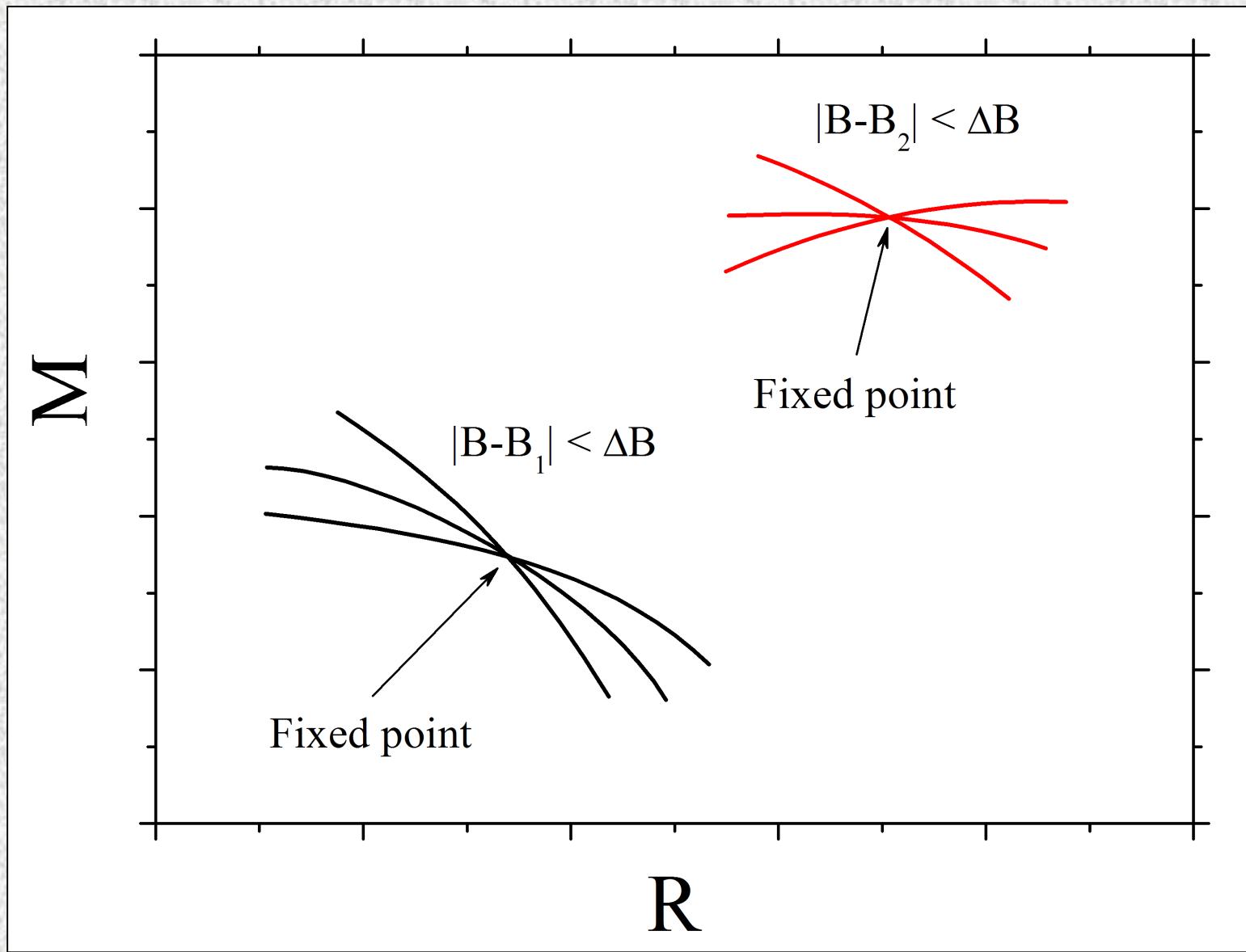


$$\left(\frac{\partial \mu}{\partial \rho_c}\right)_x \left[\rho - \frac{\alpha}{1+\alpha} + \frac{x}{2} \frac{d\rho}{dx} \right] \frac{d\rho}{dx} = x^2 \left(\frac{\partial \rho}{\partial \rho_c}\right)_x \left[\rho - \frac{\alpha}{1+\alpha} + \frac{x}{2} \frac{d\rho}{dx} \left(\rho - \frac{\mu}{x^3} \right) \right]$$

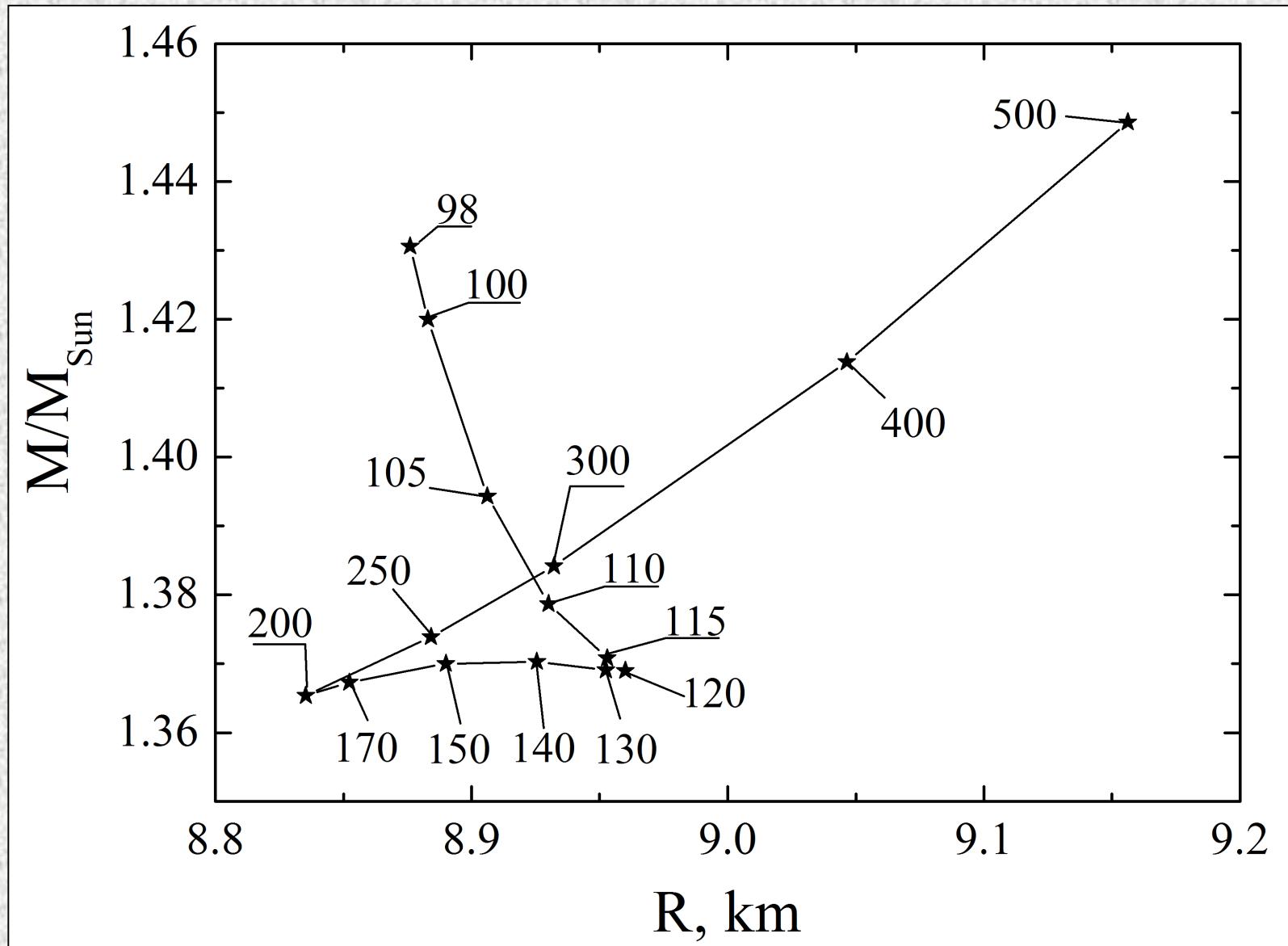


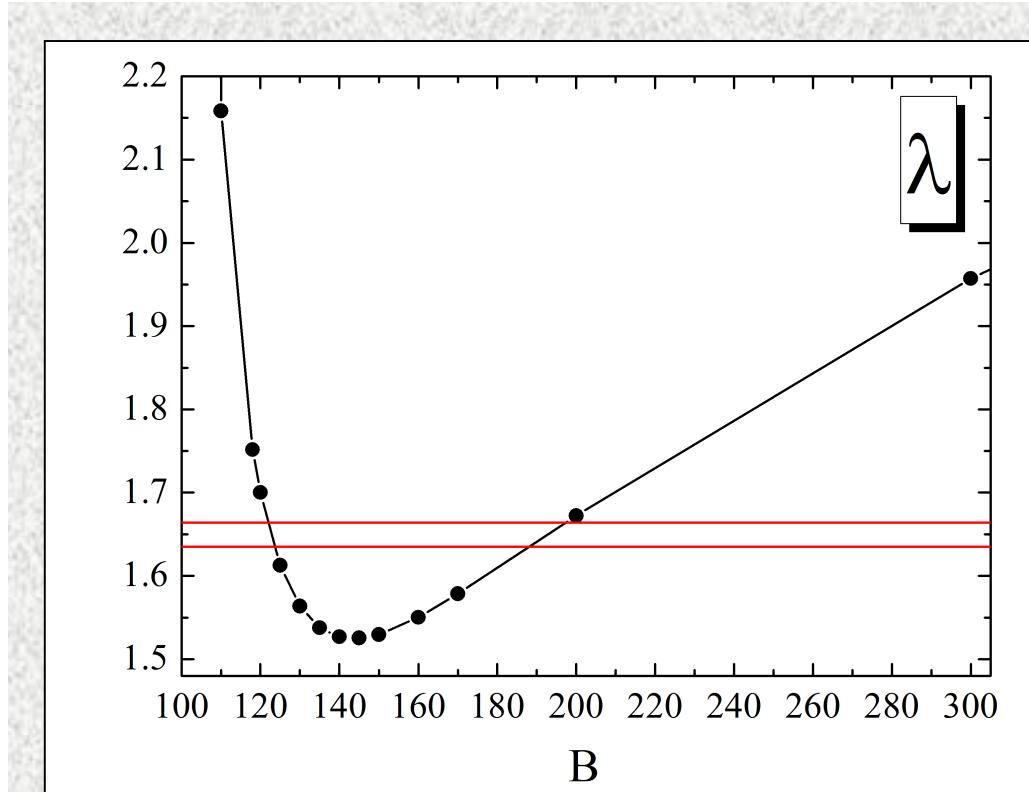
$$u_* = \left[v_*^2 (3 + \alpha) + v_* (3 - \alpha) - 6\alpha \right] \frac{1 + (1 + \alpha)(\rho_* - 1)}{4\alpha^2 (\rho_* - 1)(1 - v_*)(3 - v_*)}$$

Global structure of fixed points?



Global structure of fixed points

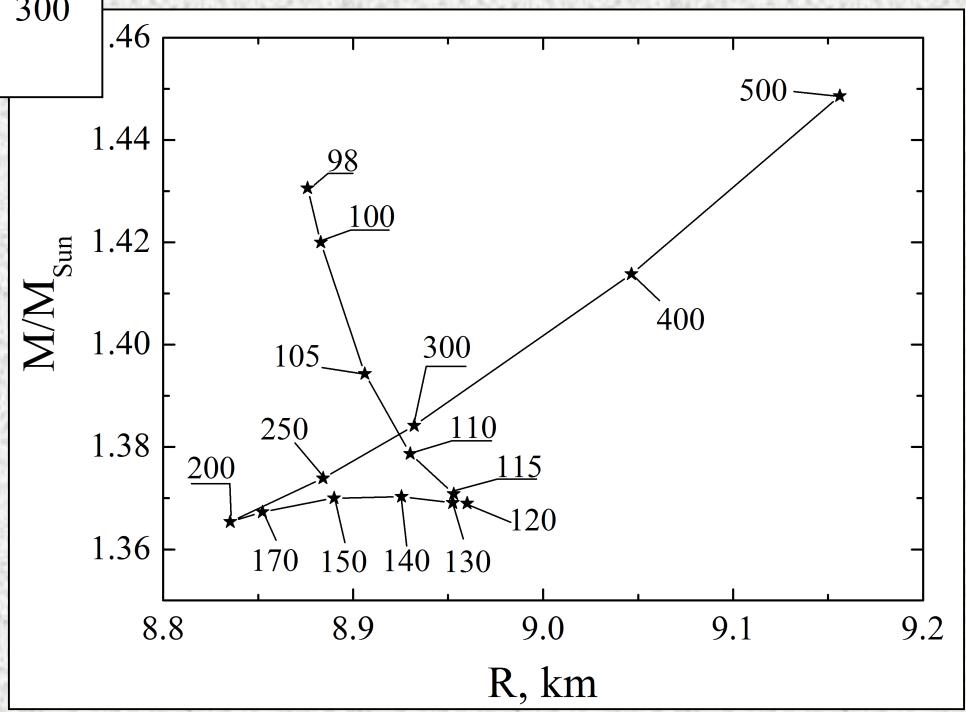




$$\lambda = \frac{u(3-v)[(1+v)^2(3-v) + 8(1-v^2)\alpha - (3-v)^2\alpha^2]}{(1+v)(7v^2-6v+3) + 8(1-v^2)(3-v)\alpha - (3-v)^3\alpha^2}$$

$$\lambda \equiv \frac{n_2}{n_1}$$

The condition that a small shift from a fixed point bring us to another fixed point



Open questions:

Other (non-linear) EOS as the solution of the main equation?

Other topology of fixed-points because different envelope?

