

Equation of state effects in core-collapse supernovae and neutrino-driven winds

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in collaboration with:

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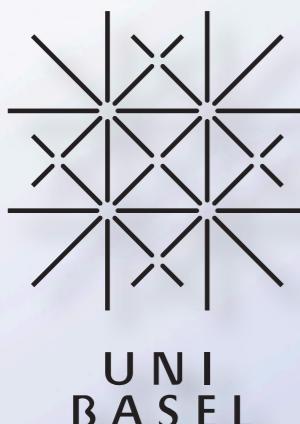
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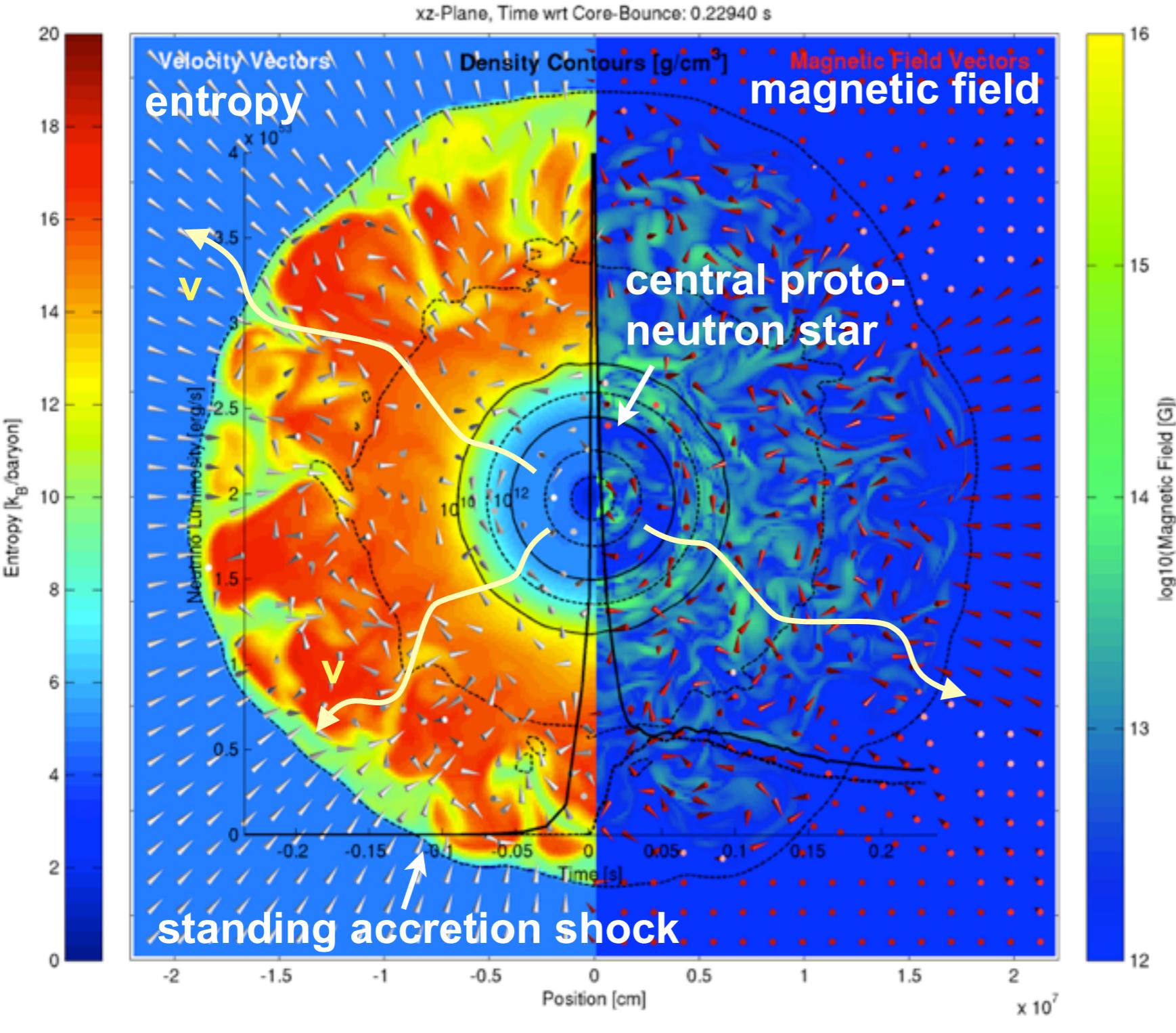


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Outline:

- 1.) introduction: supernovae, equation of state (EOS)
- 2.) EOS constraints: neutron stars, symmetry energy
- 3.) supernova simulations
- 4.) EOS effect on nucleosynthesis conditions in neutrino-driven winds
- 5.) summary

Neutrino-driven supernova mechanism



- snapshot of 3D-simulation from M. Liebendörfer
- still many open questions for core-collapse supernova explosion mechanism
- role of the EOS?

Supernova EOS – introduction

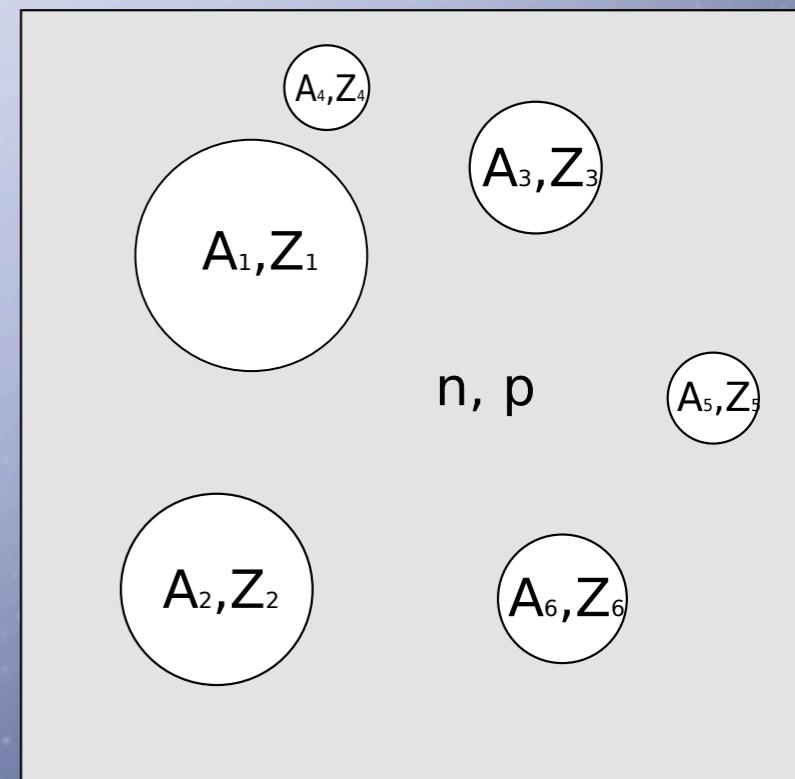
- EOS provides the nuclear physics input for astrophysical simulations: thermodynamic quantities and nuclear composition
- commonly used EOS:
 - Shen et al. (STOS): Thomas-Fermi, relativistic TM1 interactions
 - Lattimer and Swesty (LS): non-relativistic liquid drop
- plenty of EOSs for cold neutron stars
- challenge of the supernova EOS:
 - finite temperature: $T = 0 - 100 \text{ MeV}$
 - no weak equilibrium: $Y_e = 0 - 0.6$
 - wide density range: $\rho = 10^4 - 10^{15} \text{ g/cm}^3$
 - EOS in tabular form, ~ 1 million points in (T, Y_e, ρ)
- SN EOS: multi-purpose EOS, e.g., mergers
- effect on SN dynamics and nucleosynthesis?



EOS model: excluded volume NSE with interactions

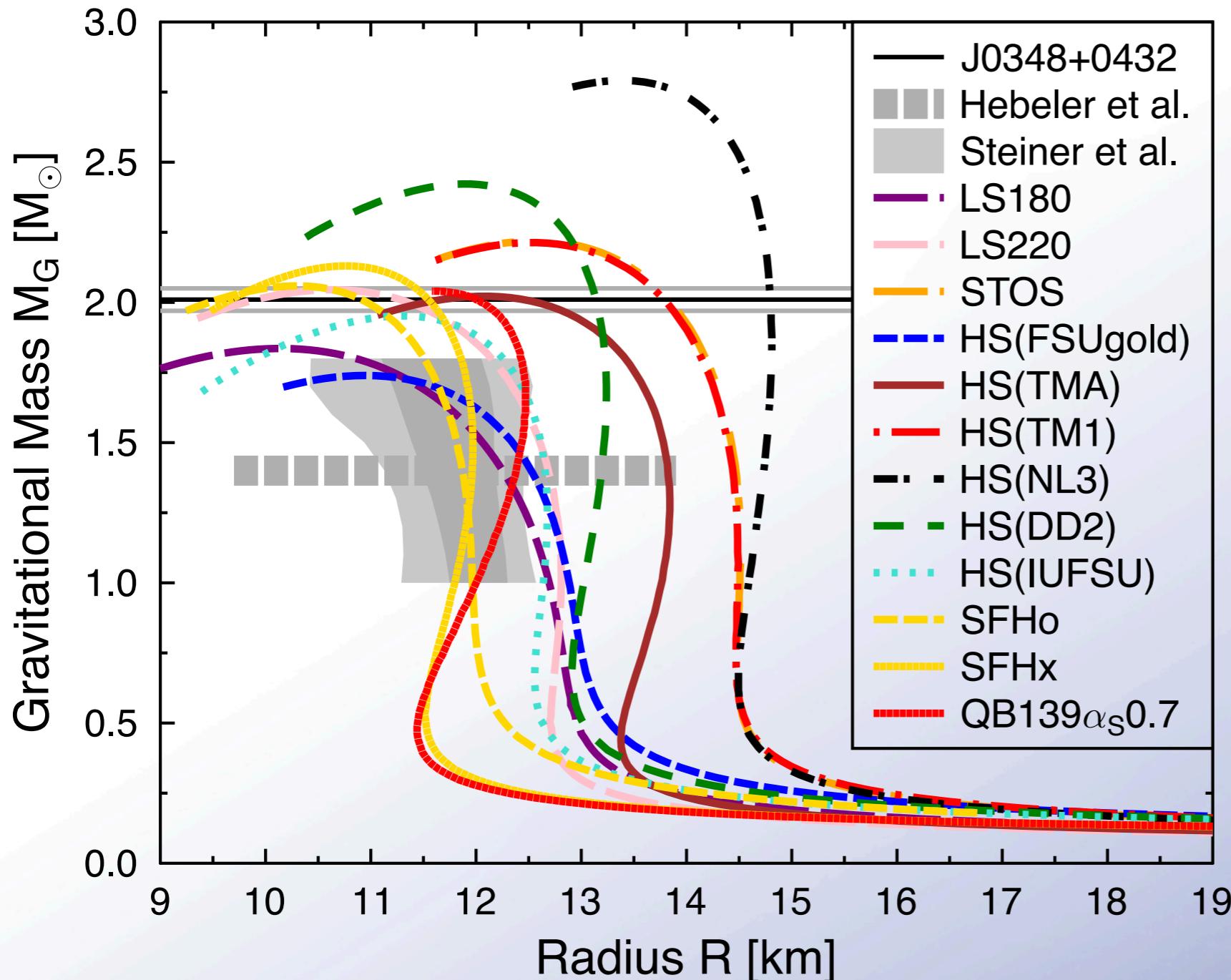
MH, J. Schaffner-Bielich; NPA 837 (2010) (HS)

- chemical mixture of nuclei and interacting nucleons in nuclear statistical equilibrium
- subsaturation densities: finite temperature generalization of outer crust EOS
 - nuclear mass tables, medium effects, Coulomb energies, excited states, excluded volume
- supersaturation densities: relativistic mean-field (RMF)
- smooth and continuous change of composition and thermodynamic quantities
- eight EOS tables for different RMF interactions:
NL3, TM1, TMA, FSUgold, DD2, SHFo, SHFx, IUFSU
<http://phys-merger.physik.unibas.ch/~hempel/eos.html>
<http://compose.obspm.fr/>



EOS constraints

EOS constraints – mass-radius relation



- PSR J0348+0432:
Antoniadis et al.
Science 2013
- Steiner et al. ApJ 2010,
Steiner et al. ApJ 2013:
bayesian analysis of NS
observations
- similar results from
Chiral EFT (Hebeler et
al. 2010)

T. Fischer, MH, et al. arXiv1307.6190

Symmetry energy

- expansion around the saturation point of nuclear matter

$$E(x, \beta) \simeq E(x, 0) + \beta^2 E_{\text{sym}}(x),$$

$$\beta = 1 - 2Y_p, \quad x = \frac{n_B}{n_0} - 1$$

$$E(x, 0) \simeq -B_0 + \frac{1}{18} K x^2 + \frac{1}{162} K' x^3 + \dots,$$

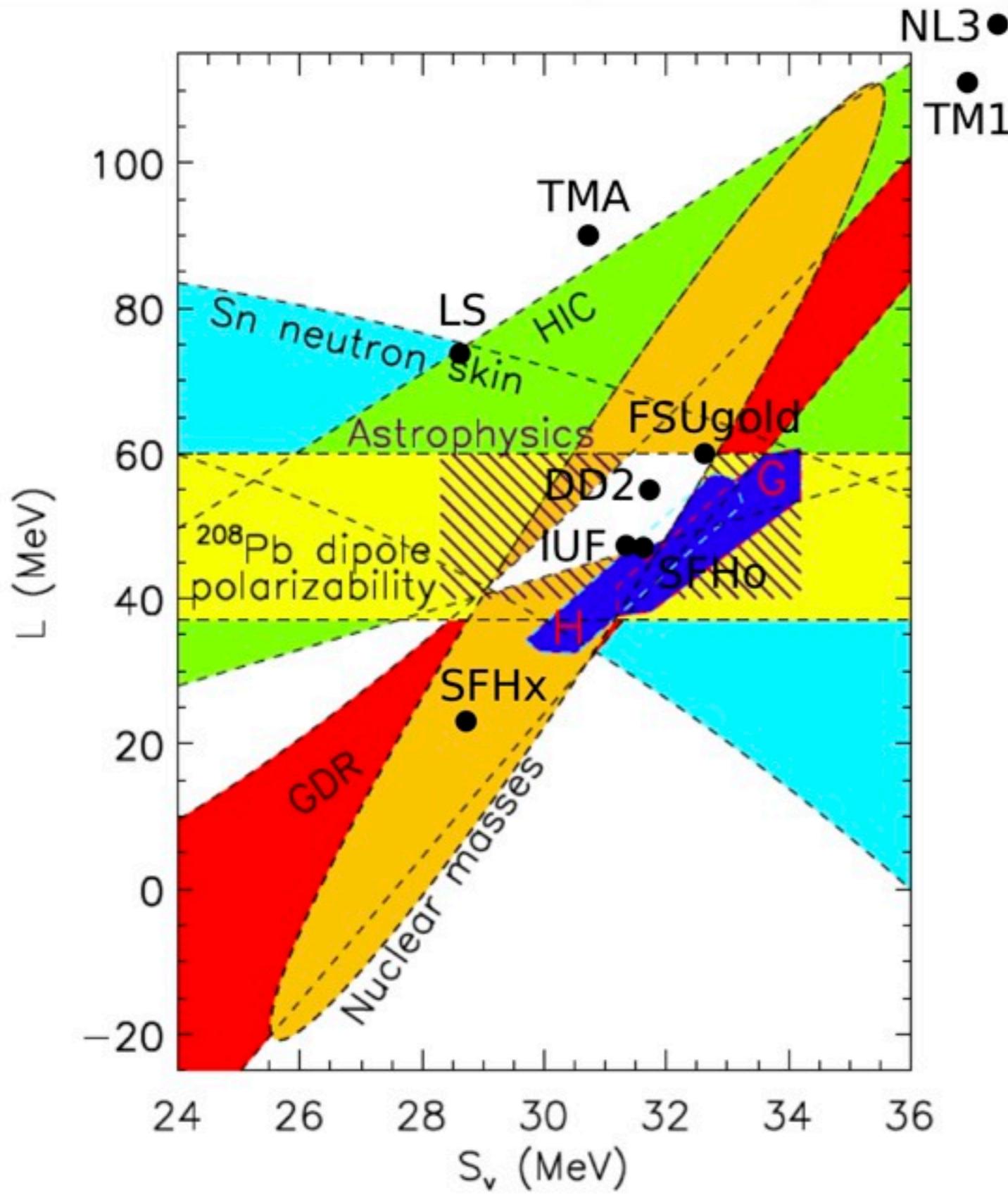
$$E_{\text{sym}}(x) \simeq J + \frac{1}{3} L x + \frac{1}{18} K_{\text{sym}} x^2 + \dots$$

- binding energy B_0 , incompressibility K , skewness K'
- symmetry energy J (S_0), slope parameter L , symmetry incompressibility K_{sym}

- L determines pressure at saturation: $p = \beta^2 n_B L / 3$
- L correlated with neutron skin thickness of heavy nuclei
- L correlated with neutron star radii
- E_{sym} determines Y_e : $\beta = \mu_e / 4E_{\text{sym}}$

EOS constraints – symmetry energy

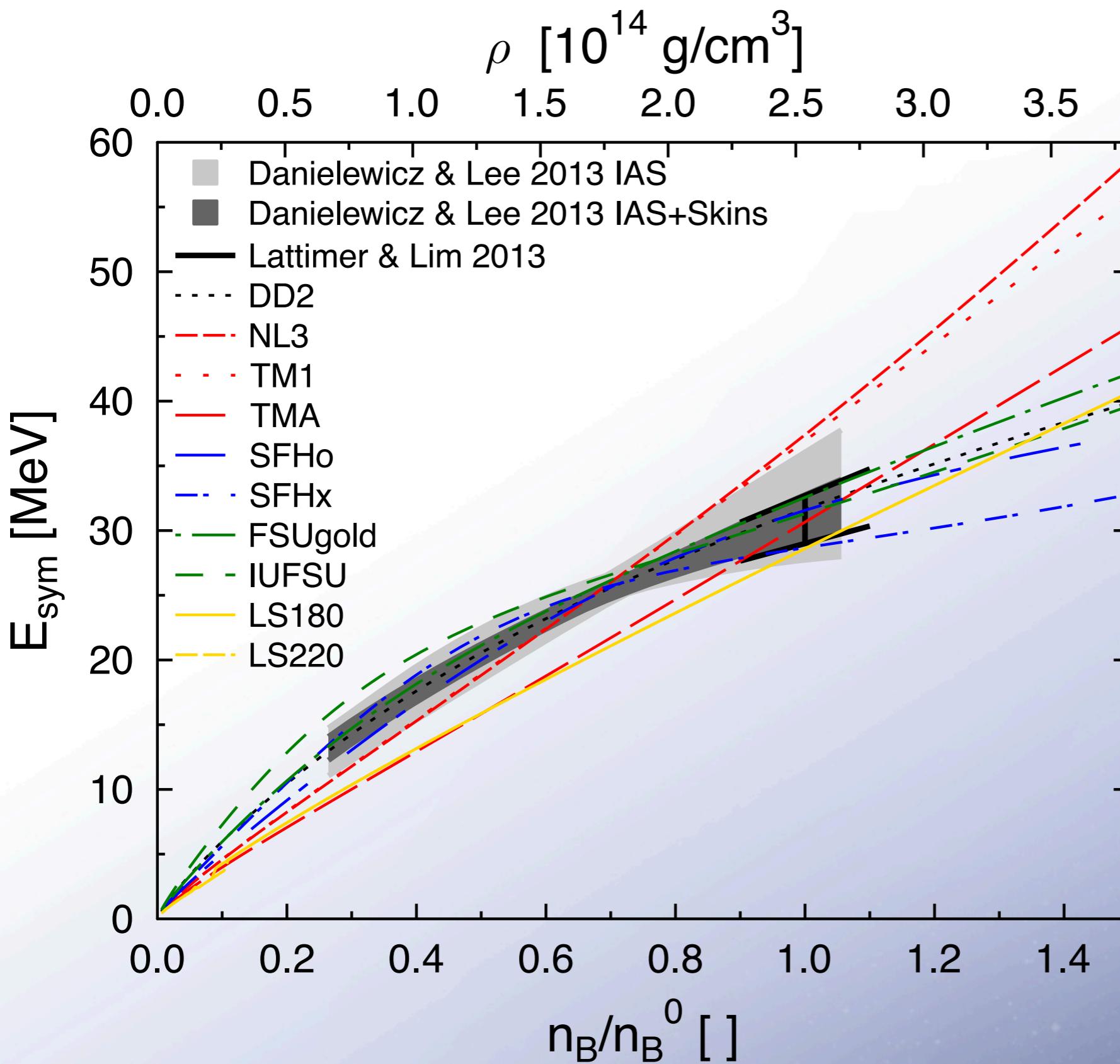
[Lattimer & Lim, ApJ 771, 51 (2013)]



- convergence of observational, experimental and theoretical constraints
- standard non-linear RMF in disagreement (TM1,NL3,TMA)
- SFHo, IUF and DD2 perform well
- G: Gandolfi et al. 2012: quantum Monte-Carlo
- H: Hebeler et al. 2010: Chiral EFT, neutron matter

EOS constraints – symmetry energy

based on:
[Danielewicz & Lee, arXiv:1307.4130]



- derived from binding energies of isobaric analog states
- standard non-linear RMF (NL3, TM1, TMA) in disagreement
- LS EOS extremely low
- DD2: perfect match

$n_0 \sim 0.16 \text{ fm}^{-3}$
 $\sim 2 \times 10^{14} \text{ g/cm}^3$



Supernova simulations

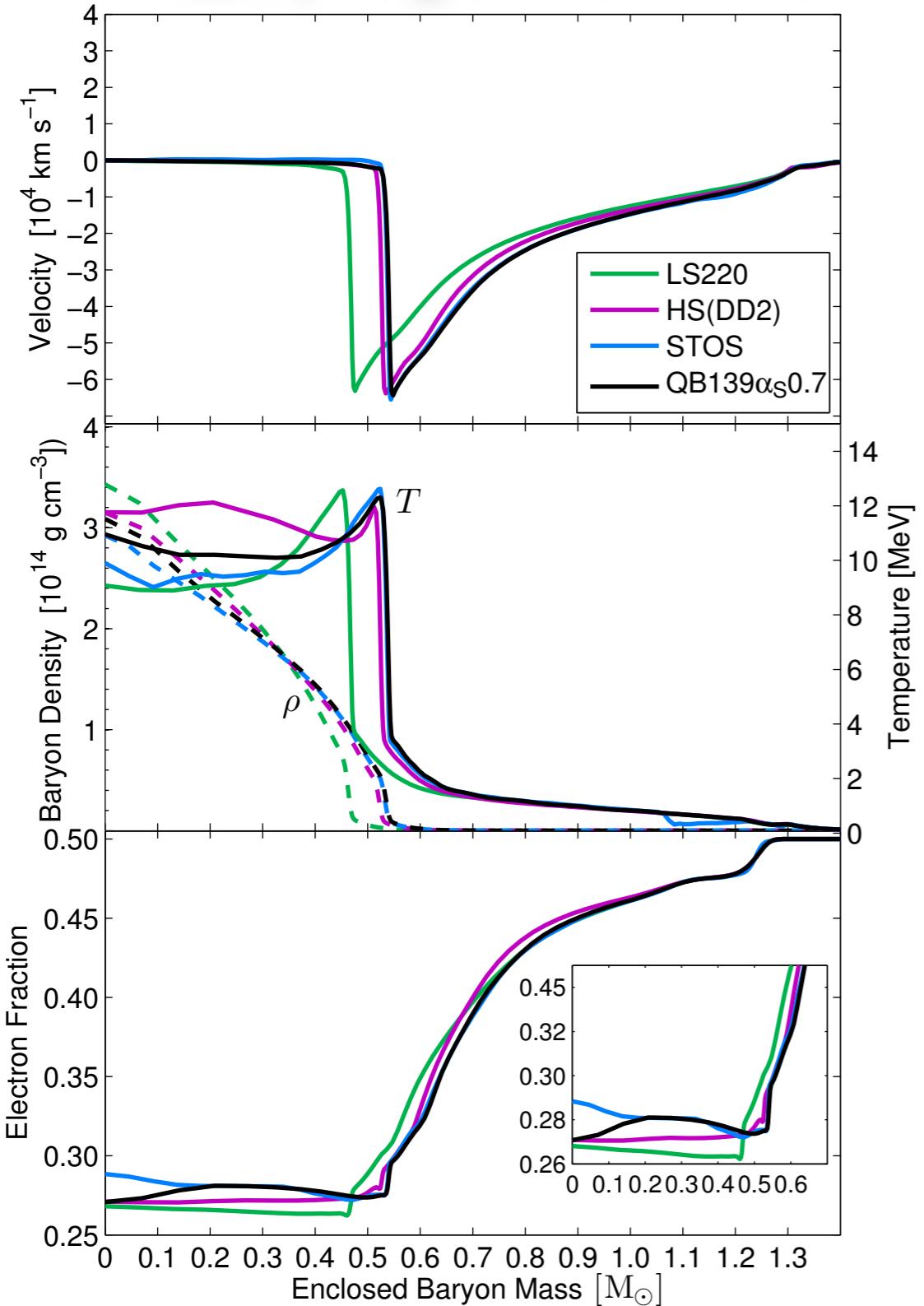


Supernova simulations – setup

MH, T. Fischer, J. Schaffner-Bielich, M. Liebendörfer, ApJ 748, 70 (2012),
A. Steiner, MH, T. Fischer; ApJ 774, 17 (2013),
T. Fischer, MH, et al. arXiv1307.6190

- simulations by Tobias Fischer, University of Wroclaw
 - general relativistic radiation hydrodynamics in spherical symmetry
 - three flavor Boltzmann neutrino transport
- $11.2 \text{ M}_{\text{sun}}$ progenitor of Woosley et al. RMP 74 (2002)
 - regular core-collapse supernovae
 - comparison of LS220, STOS and HS(DD2)

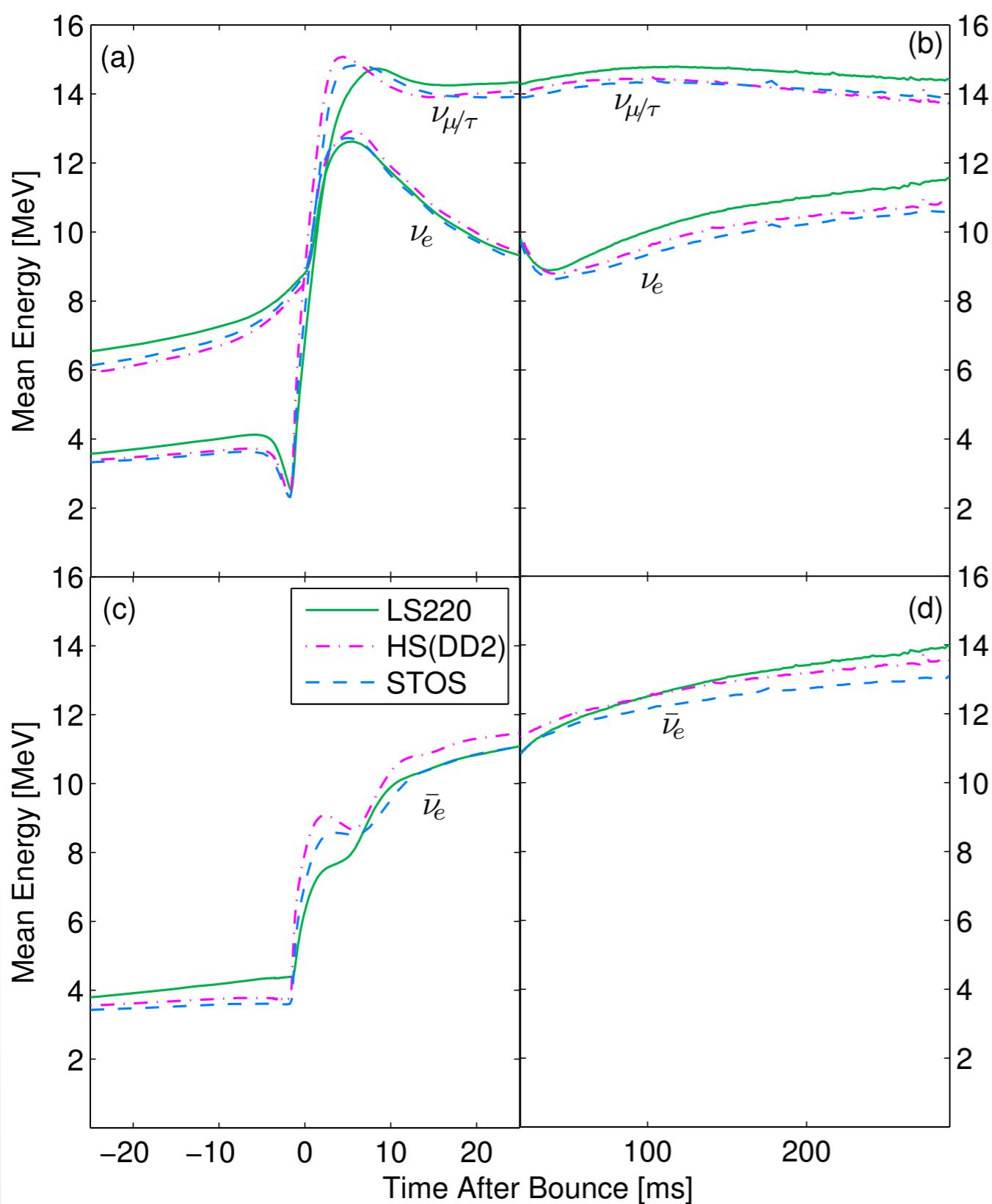
$11 M_{\odot}$ progenitor – bounce



- highest densities, lowest bounce mass, lowest Y_e for LS220 \leftrightarrow lowest J
- DD2 between LS220 and STOS

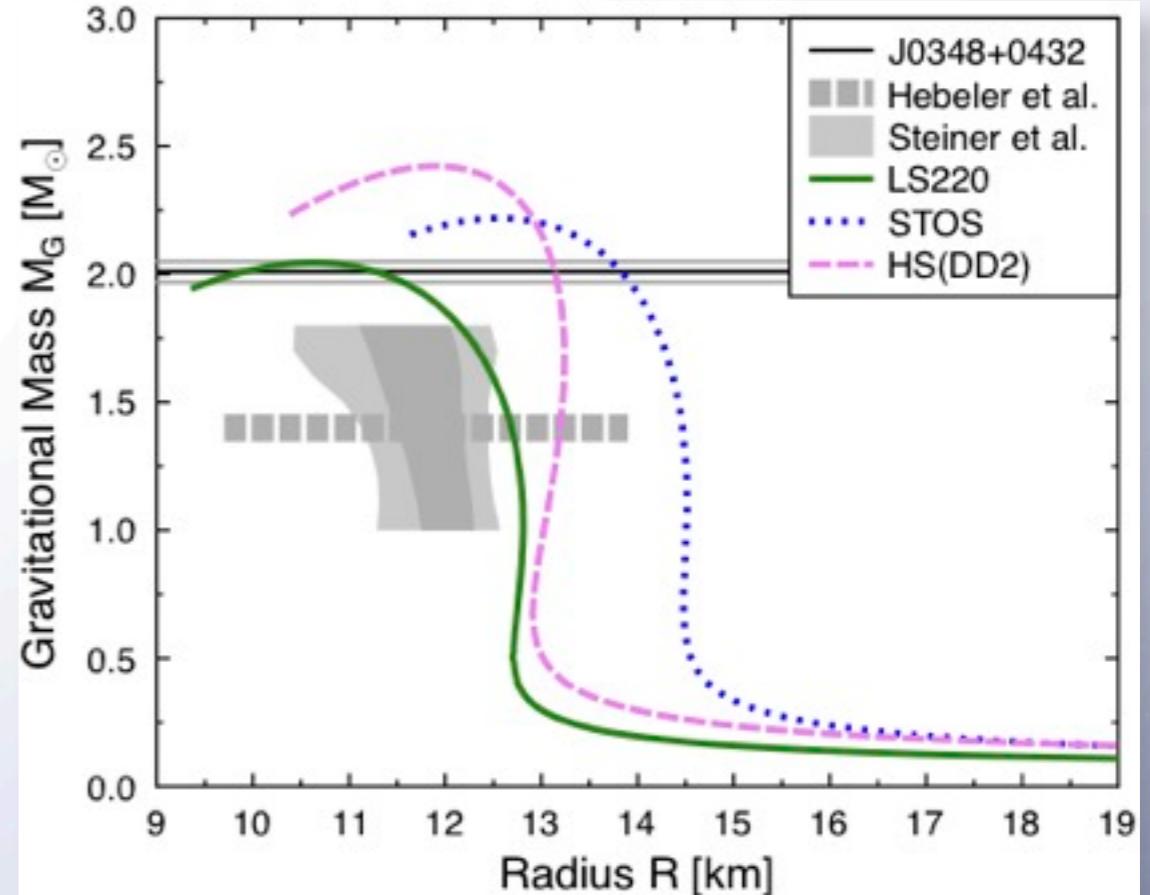
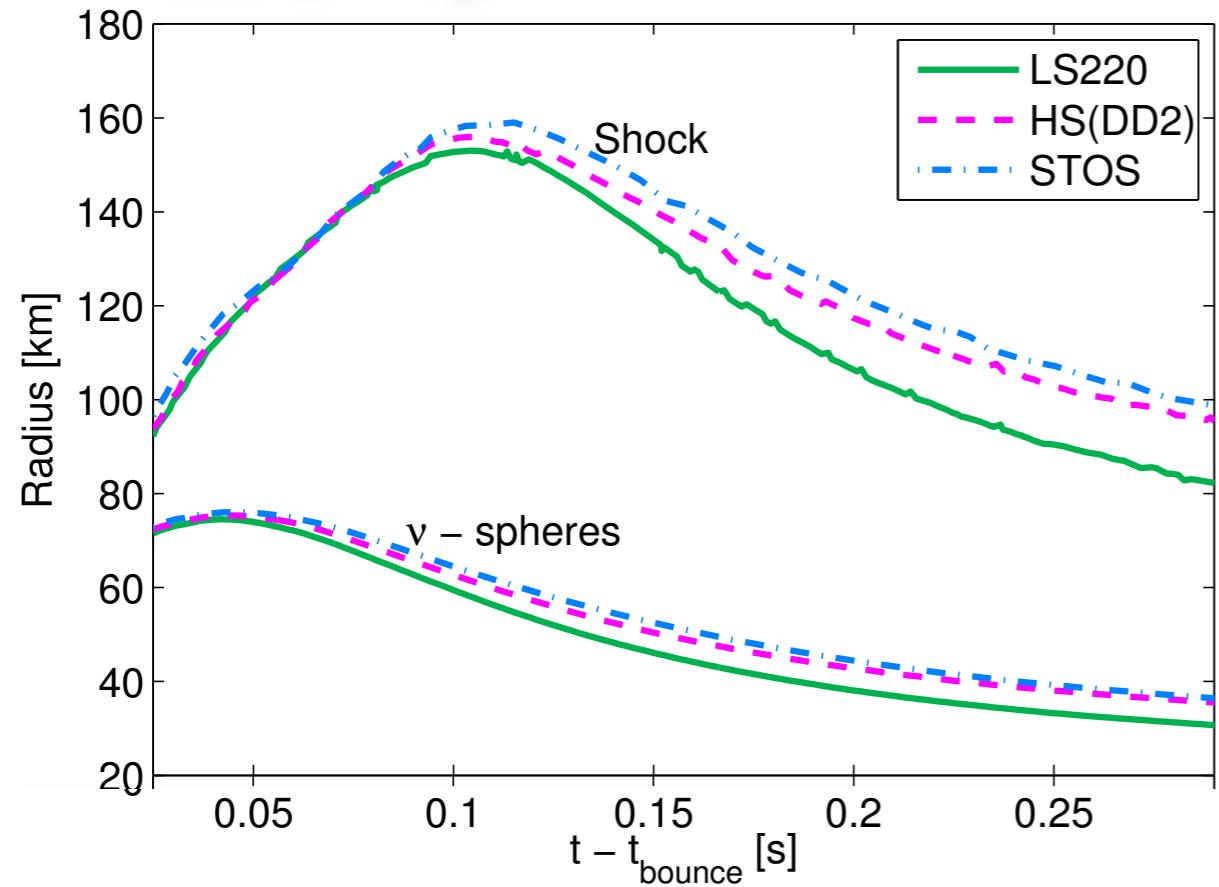
EOS	J [MeV]
LS220	28.6
HS(DD2)	31.7
STOS	36.9

$11 M_{\text{sun}}$ progenitor – Neutrinos



- ordering of mean energies:
 $\text{LS220} > \text{HS(DD2)} > \text{STOS}$

$11 M_{\text{sun}}$ progenitor – Shock evolution



- higher mean energies of LS220 as a result of faster proto-neutron star contraction
- DD2 between the more extreme cases of STOS and LS
- radius evolution in agreement with TOV solutions
- small differences, but shown to be amplified in Multi-D (Suwa et al. 2013, Marek et al, 2009, Couch 2013)

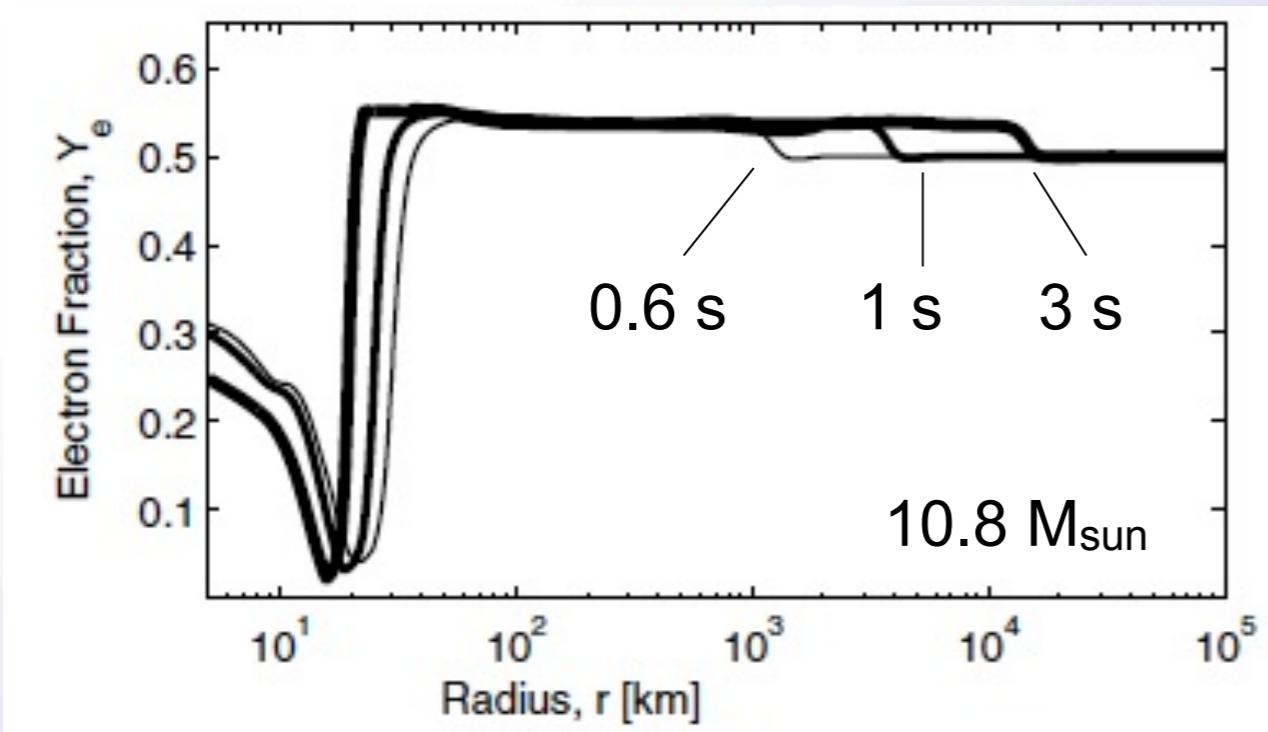
STOS too pessimistic, LS220 too optimistic for explosion models

Neutrino-driven winds

Neutrino-driven winds

- after onset of explosion: neutrino emission from newly born proto neutron star
 - neutrino-driven wind: energy deposition by neutrinos leads to emission of low-density, high-entropy matter from proto-neutron star surface
 - candidate site for r-process nucleosynthesis
-
- previous long-term core-collapse supernova simulations by Fischer et al (2010), Hüdepohl et al. (2010):

[Fischer et al. A&A 517 (2010)]



- the neutrino-driven wind is generally proton rich
- allows vp-process
(C. Fröhlich et al. 2006, Pruet et al. 2006, Wanajo et al. 2006, Arcones et al. 2011, Arcones & Thielemann 2012, ...)

Estimate for Y_e

- Qian & Woosley, ApJ 471 (1996):

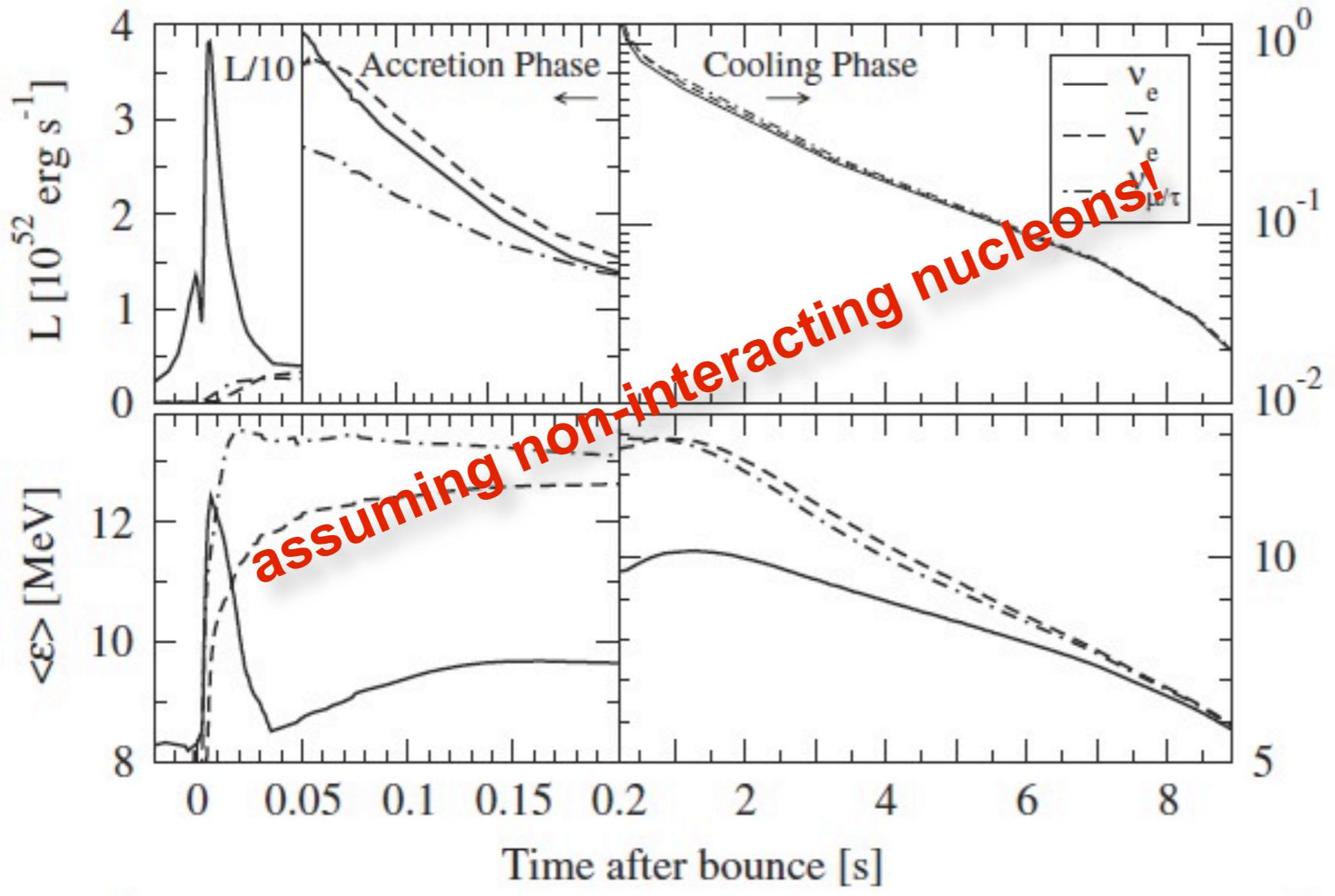
$$Y_e \simeq \left(1 + \frac{L_{\bar{\nu}_e} \langle \epsilon_{\bar{\nu}_e} \rangle - 2Q + \frac{1.2 Q^2}{\langle \epsilon_{\bar{\nu}_e} \rangle}}{L_{\nu_e} \langle \epsilon_{\nu_e} \rangle + 2Q + \frac{1.2 Q^2}{\langle \epsilon_{\nu_e} \rangle}} \right)^{-1}$$

- $Q = m_n - m_p = 1.3 \text{ MeV}$
- neglects neutrino emission, i.e. electron and positron captures
- for similar luminosities:

$$Y_e < 0.5 \Leftrightarrow E_{\bar{\nu}_e} - E_{\nu_e} > 4Q$$

Neutrino properties in the “standard” wind

[Hüdepohl et al. PRL 104 (2010)]



- similar luminosities
- $E_{\nu_{\mu/\tau}} \sim E_{\bar{\nu}_e} > E_{\nu_e}$
- $E_{\bar{\nu}_e} - E_{\nu_e} < 4Q \rightarrow Y_e > 0.5$

Mean-field potentials in charged-current rates

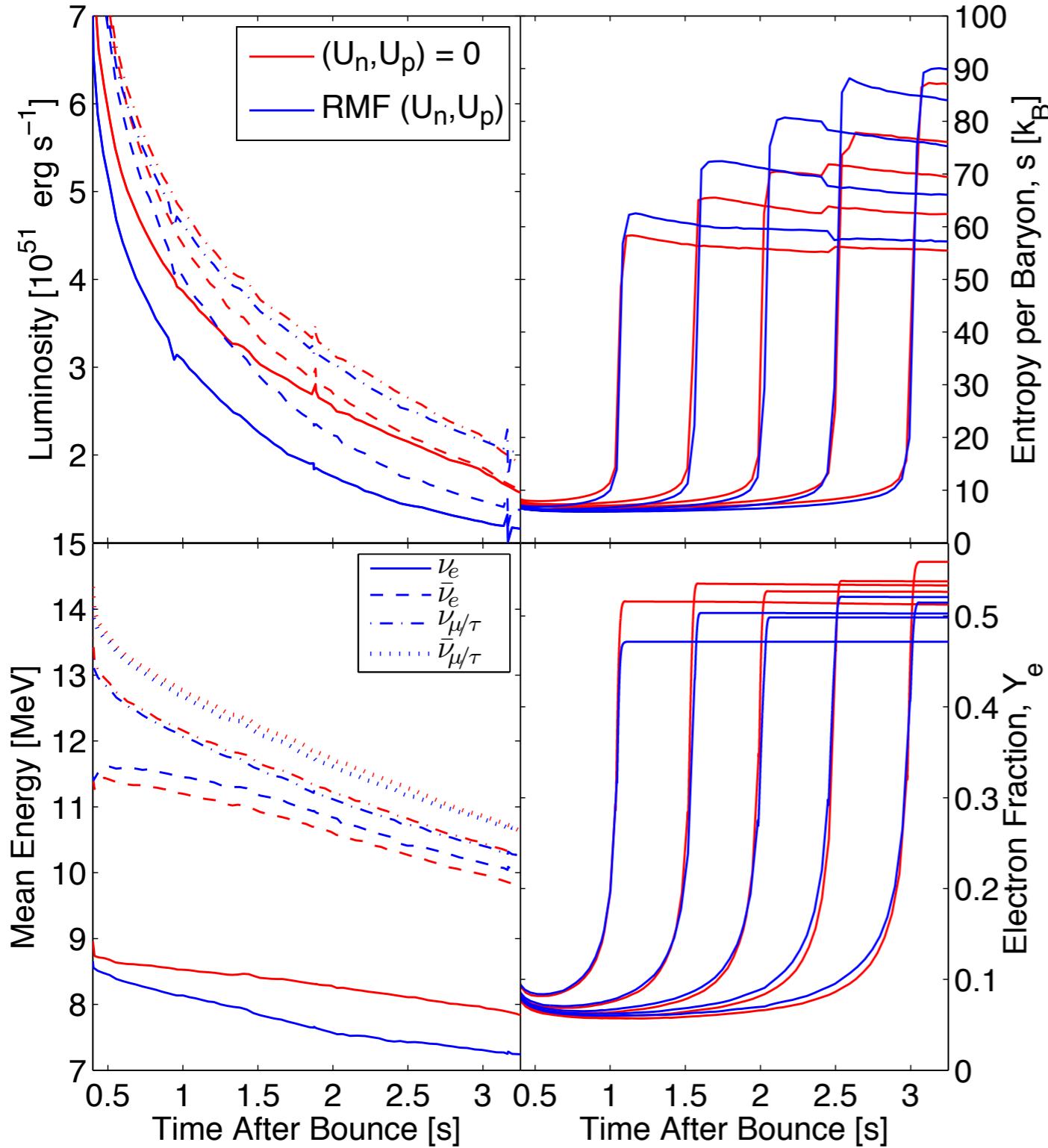
- Bruenn '85: charged-current rates based on non-interacting nucleons
- improved charged-current rates (with mean-field effects):
 - Reddy, Prakash & Lattimer, PRD58 (1998)
 - Reddy, Prakash, Lattimer & Pons, PRC59 (1999)
- G. Martínez-Pinedo et al., PRL109 (2012), Roberts & Reddy, PRC86 (2012): crucial for late neutrino spectra
- nucleons in self-generated mean-field potentials U_i
- energy conservation for a generic mean-field model e.g.: $e + p \rightarrow n + \nu_e$

$$\begin{aligned} E_p + E_e &= E_n + E_{\nu_e} \\ \sqrt{p_p^2 + m_p^{*2}} + U_p + \sqrt{p_e^2 + m_e^2} &= \sqrt{p_n^2 + m_n^{*2}} + U_n + E_{\nu_e} \end{aligned}$$

- neutron-rich conditions: $\Delta U = U_n - U_p > 0 \rightarrow$ reduces neutrino energies
- only relevant at high densities

Effects of mean-field potentials on the neutrino-driven wind

[Martínez-Pinedo et al., PRL 109 (2012)]



- E_{νe} decrease, E_{ν̄e} increased, difference increased
- mean-field effects lead to slightly neutron-rich wind
- same conclusions by Roberts et al. 2012, different EOS: IUFSU, lowest Y_e of 0.43

definition of nucleon mean-field potentials

What sets ΔU ?

- for the eight relativistic SN EOS available:

$$\begin{aligned} E &= E^{\text{kin}} + E^{\text{int}} \\ E^{\text{int}} &= E^{\text{int}}(T, n_B, Y_p, \Sigma_S, \Sigma_V^i) \end{aligned}$$

- Σ_S : scalar self energy
- Σ_V : vector self energy

$$\Delta U \simeq 4(1 - 2Y_e)E_{\text{sym}}^{\text{int}}$$

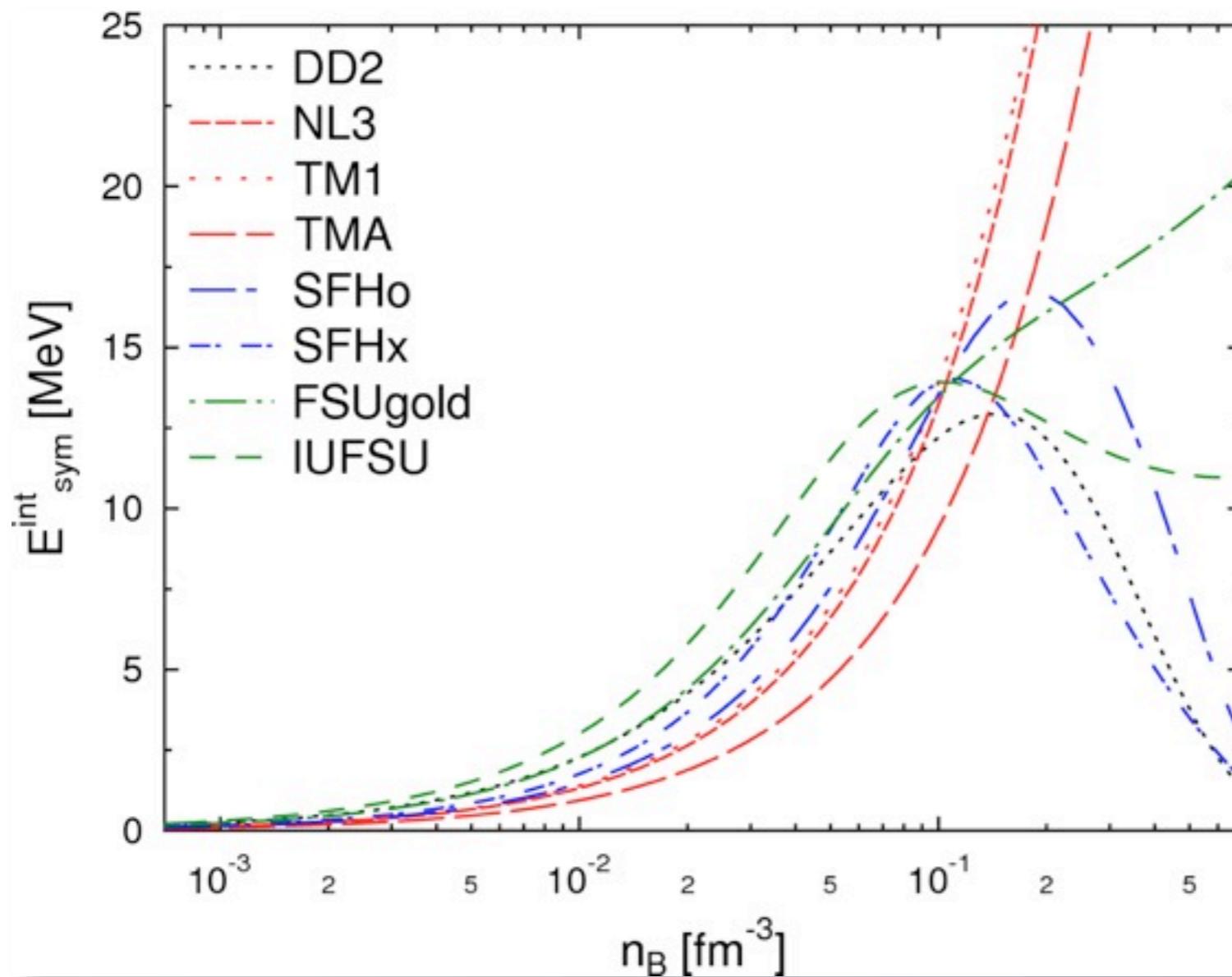
$$E_{\text{sym}}^{\text{int}}(n_B) = \frac{1}{8n_B} \left. \frac{d}{dY_p} \right|_{n_B, T=0} \frac{\partial E^{\text{int}}}{\partial Y_p}$$

- interaction part of the symmetry energy

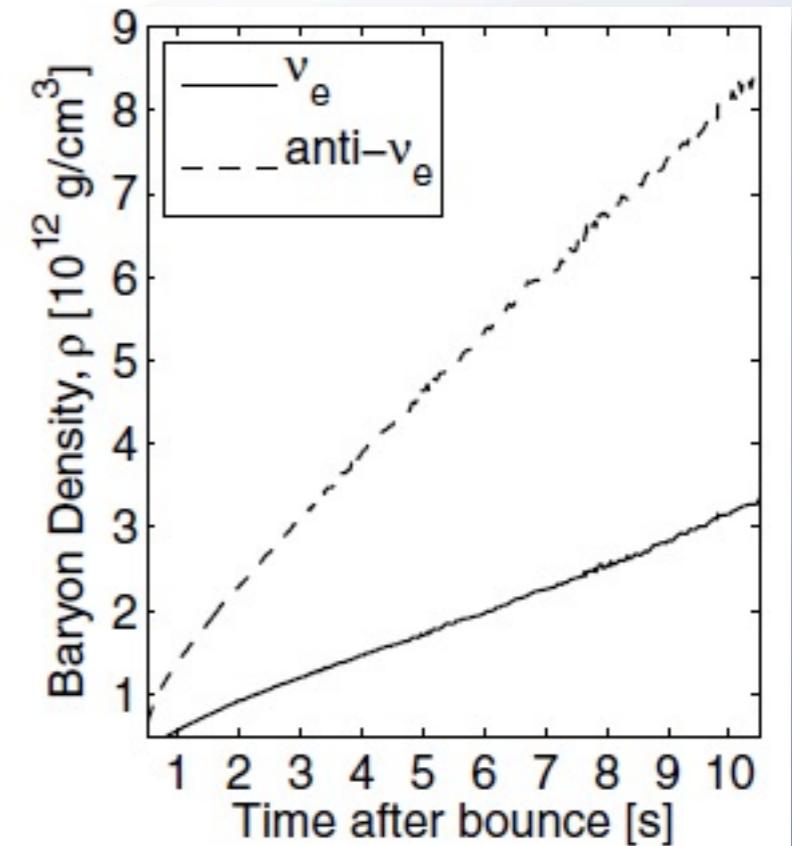
$$E_{\text{sym}} = E_{\text{sym}}^{\text{int}} + E_{\text{sym}}^{\text{kin}}$$

Interaction part of the symmetry energy

$n_0 \sim 0.16 \text{ fm}^{-3}$
 $\sim 2 \times 10^{14} \text{ g/cm}^3$



neutrinosphere

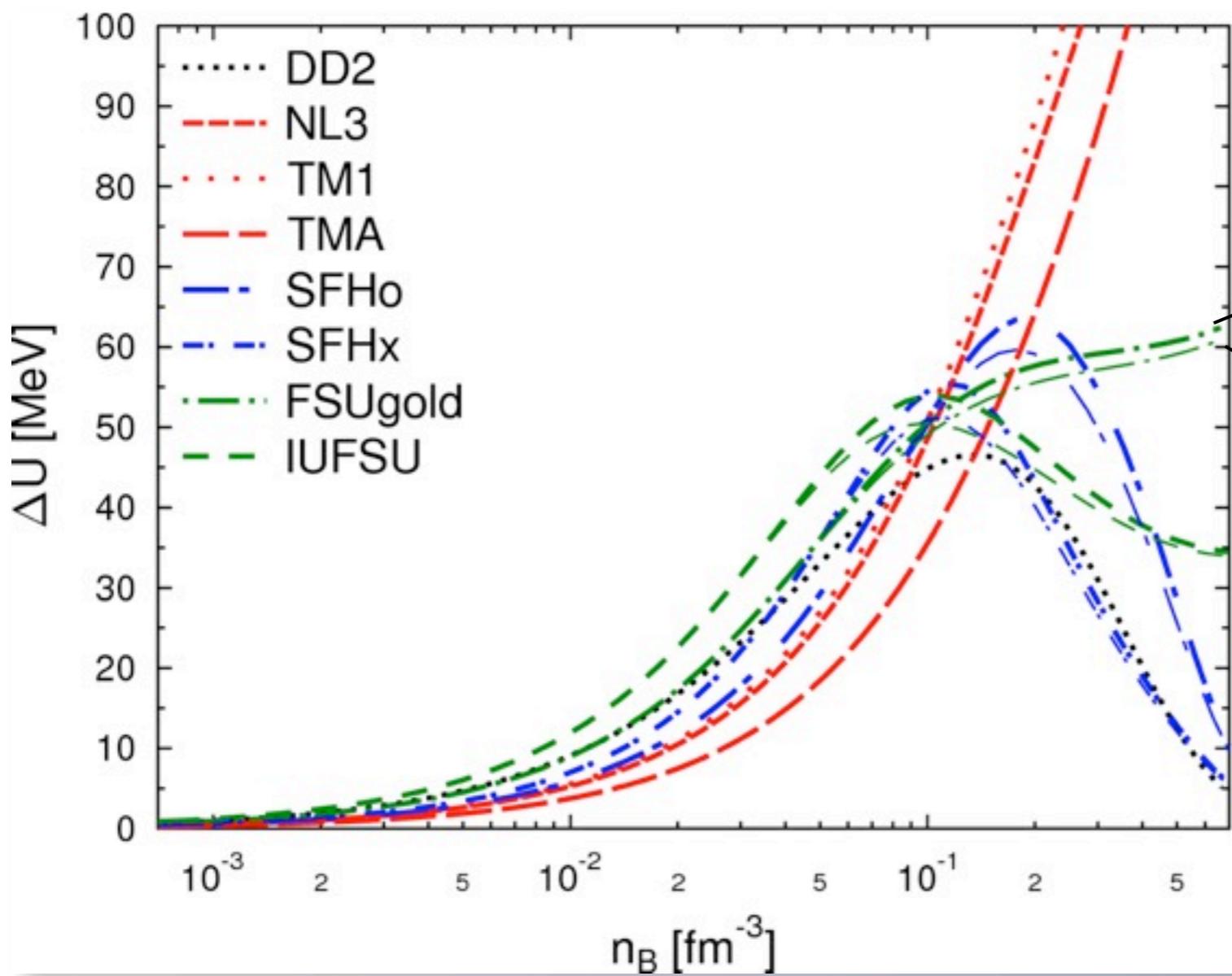


[Fischer et al. A&A 517 (2010)]

- even at low densities substantial differences
- diverging for $n_B > 0.1 \text{ fm}^{-3}$
- compensation effect: higher E_{sym} \rightarrow higher Y_e

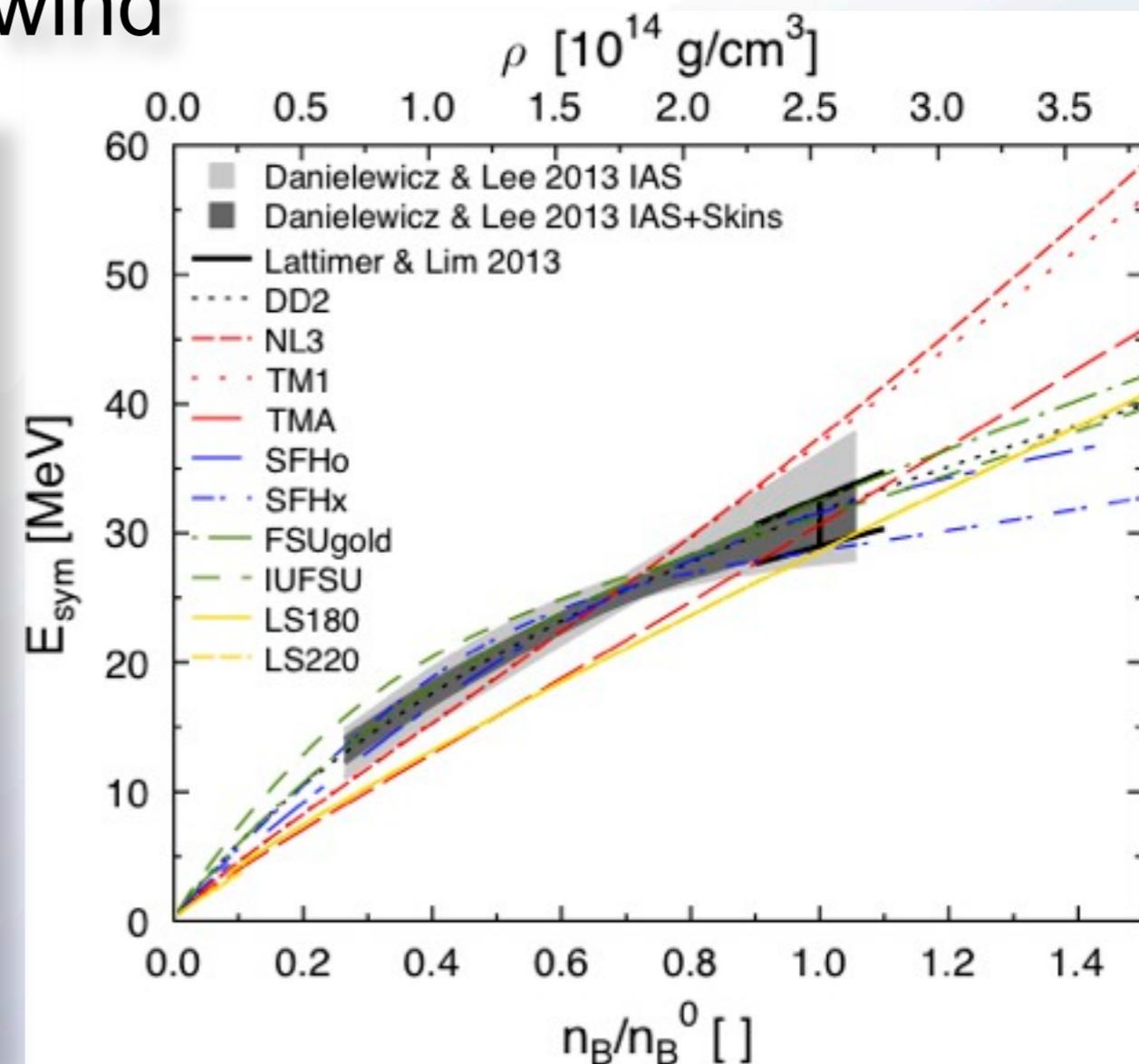
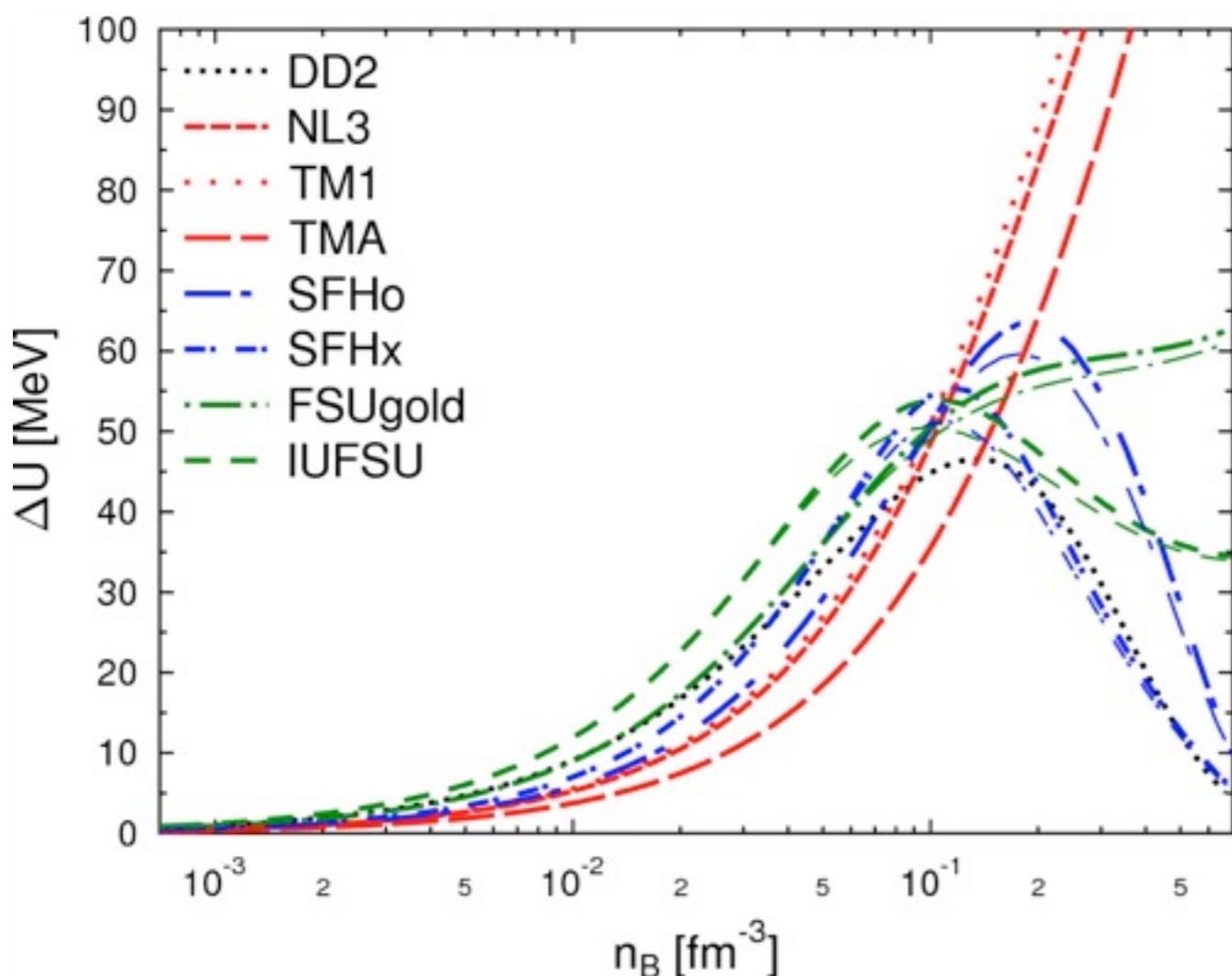
$$\Delta U \simeq 4(1 - 2Y_e)E_{\text{sym}}^{\text{int}}$$

ΔU in beta-equilibrium



- beta-eq: upper limit for ΔU
 - compensation by Y_e less important than differences due to $E_{\text{sym}}^{\text{int}}$
 - exact
 - approx. by $E_{\text{sym}}^{\text{int}}$
- $$\Delta U \simeq 4(1 - 2Y_e)E_{\text{sym}}^{\text{int}}$$
- approximation accurate up to 10%

Expectations for neutrino-driven wind

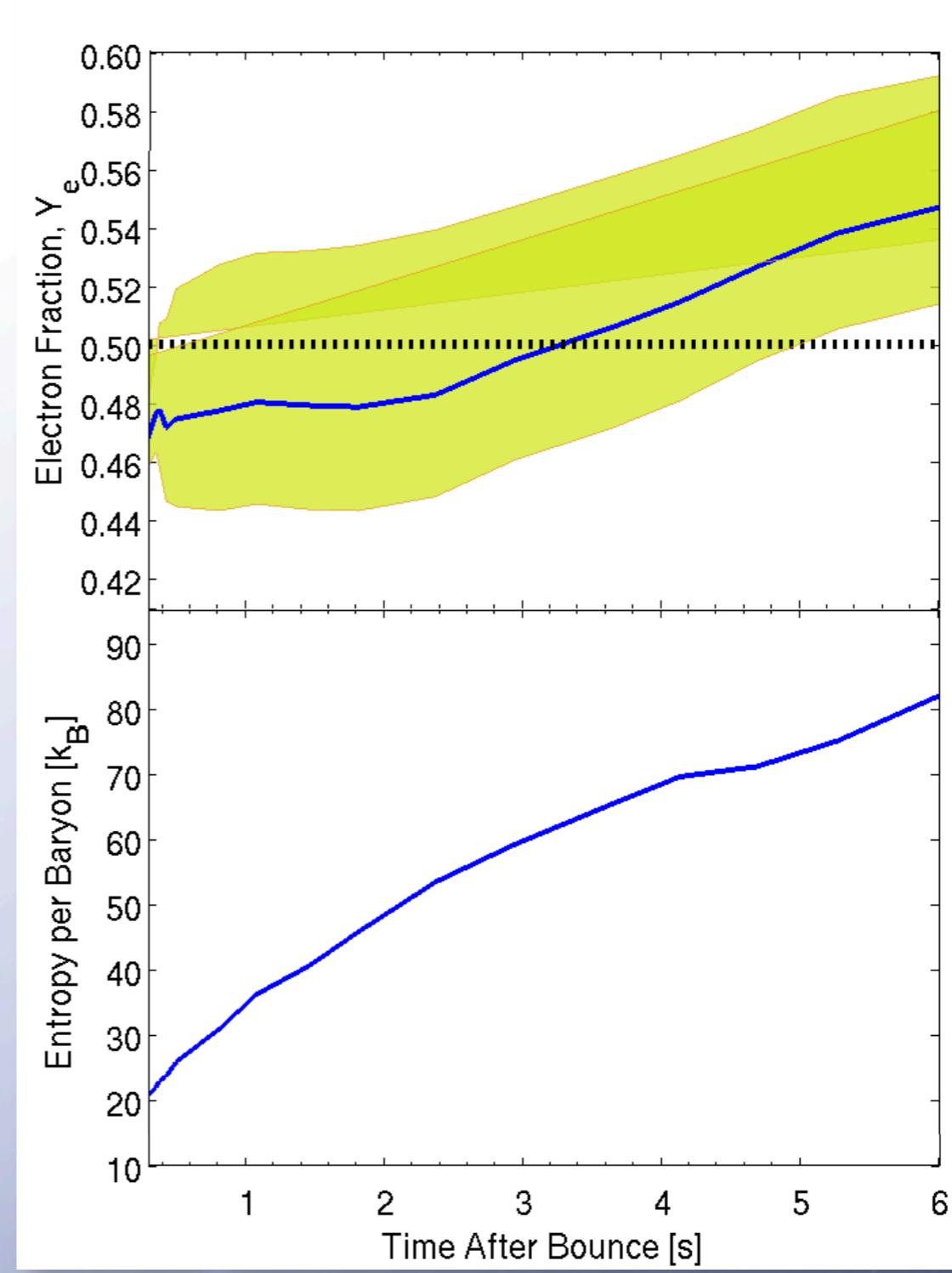
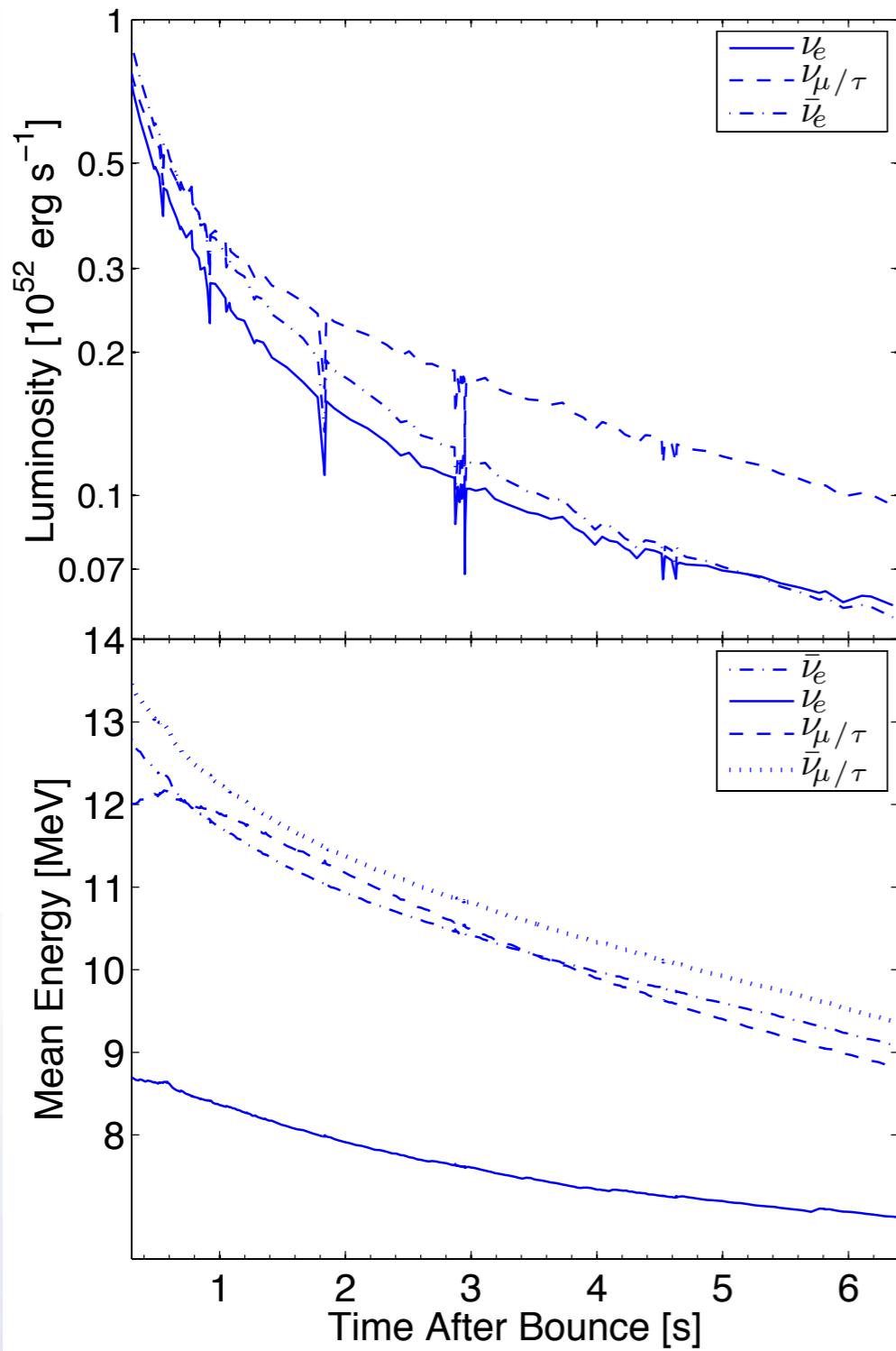


- tested so far: TM1 → lowest Y_e of 0.47 at $\sim 1 \text{ s pb}$
- most realistic EOS: DD2 with higher ΔU at low densities → lower Y_e ?

- scattering vs. charge current rates?
- uncertainties in rates (e.g. Bremsstrahlung)?
- effects of correlations, light nuclei?
- momentum-dependent interactions?

Wind simulations for DD2 EOS

- (preliminary) results for DD2 EOS with artificial explosion by T. Fischer, L. Huther and G. Martinez-Pinedo



Conclusions

- EOS tables and routines for composition available for NL3, TM1, TMA, FSUgold, IUFSU DD2, SFHo, SFHx:
<http://phys-merger.physik.unibas.ch/~hempel/eos.html>

<http://compose.obspm.fr/>

CompOSE

CompStar Online
Supernovae Equations of State



- symmetry energy and neutron matter EOS at and below n_0 well constrained
- puts constraints on the dynamics of supernovae
- classical SN EOS (LS and STOS) disfavored, new EOS models indicate an intermediate behavior in SN → multi-D studies
- interaction symmetry energy influences neutrino spectra evolution
- the asymmetry at high densities determines the asymmetry of the wind ejecta
- uncertainties, implications on nucleosynthesis?